

LAPLACE TRANSFORMS

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WE DEFINE THE LAPLACE TRANSFORM BY

$$\hat{F}(s) = \int_0^{\infty} e^{-st} f(t) dt = \mathcal{L}(f(t)).$$

WE REQUIRE $|f(t)| \leq K e^{at}$ FOR ALL t , AND $f(t)$ IS PIECEWISE CONTINUOUS. THEN, WE OBTAIN THAT $F(s)$ IS ANALYTIC FOR $\text{RE}(s) > a$.

PROPERTIES

(i) $\mathcal{L}(f_1 + f_2) = \mathcal{L}(f_1) + \mathcal{L}(f_2)$ LINEARITY

(ii) $\mathcal{L}(f'') = s^2 \mathcal{L}(f) - f'(0) - s f(0)$. $\mathcal{L}(f') = s \mathcal{L}(f) - f(0)$.

(iii) $\mathcal{L}(f(t) e^{at}) = F(s-a)$, WHERE $F(s) = \mathcal{L}(f(t))$.

(iv) $\mathcal{L}(t f(t)) = \int_0^{\infty} t f(t) e^{-st} dt = -F'(s)$, WHERE $F(s) = \mathcal{L}(f(t))$.

(v) $f(t+\tau) = f(t)$. THEN,

$$\begin{aligned} \mathcal{L}(f(t)) &= \int_0^{\infty} f(t) e^{-st} dt = \int_0^{\tau} f(t) e^{-st} dt + \int_{\tau}^{2\tau} e^{-st} f(t) dt \\ &\quad + \dots + \int_{(n-1)\tau}^{n\tau} f(t) e^{-st} dt + \dots \\ &= \int_0^{\tau} f e^{-st} dt (1 + e^{-s\tau} + e^{-2s\tau} + \dots) \end{aligned}$$

HENCE,

$$\mathcal{L}(f(t)) = \frac{\int_0^{\tau} e^{-st} f(t) dt}{1 - e^{-s\tau}}.$$

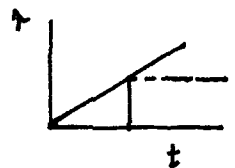
(vi) LET $u_c(t) = \begin{cases} 0 & 0 \leq t < c \\ 1 & t \geq c \end{cases}$ UNIT STEP FUNCTION.

THEN, $\mathcal{L}(u_c(t) f(t-c)) = \int_c^{\infty} e^{-st} f(t-c) dt = e^{-sc} F(s)$.

(vii) CONVOLUTION $(g * h)(t) = \int_0^t g(t-\tau) h(\tau) d\tau = f(t)$.

THEN, $\mathcal{L}[(g * h)(t)] = \hat{G}(s) \hat{H}(s)$. $\hat{G}(s) = \mathcal{L}(g)$, $\hat{H}(s) = \mathcal{L}(h)$.

PROOF $\hat{F}(s) = \int_0^{\infty} e^{-st} \left(\int_0^t g(t-\tau) h(\tau) d\tau \right) dt$



NOW INTERCHANGING ORDER OF INTEGRATION,

$$\hat{F}(s) = \int_0^{\infty} h(\tau) \int_{\tau}^{\infty} e^{-st} g(t-\tau) dt d\tau = \int_0^{\infty} h(\tau) e^{-s\tau} \left(\int_{\tau}^{\infty} e^{-s(t-\tau)} g(t-\tau) dt \right) d\tau$$

LETTING $u = t - \tau$, $du = dt$, WE GET $\hat{F}(s) = \hat{G}(s) \hat{H}(s)$

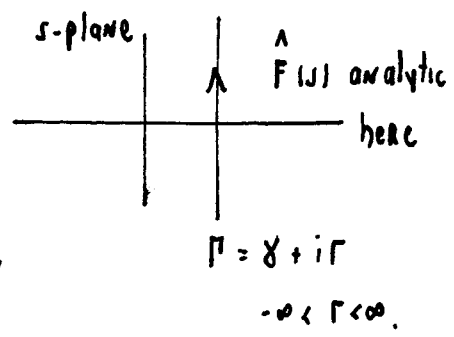
(viii) INVERSION FORMULA

LET $\hat{F}(s) = \int_0^\infty e^{-st} f(t) dt.$

SUPPOSE THAT THE INTEGRAL (CONVERGES) FOR SOME s_0 , THEN IT CONVERGES FOR ALL s WITH $RE(s) > RE(s_0)$; IN THIS RANGE $\hat{F}(s)$ IS ANALYTIC.

NOW WE CLAIM,

$f(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} \hat{F}(s) e^{st} ds$



where $\hat{F}(s)$ IS ANALYTIC FOR $RE(s) > \gamma$. HERE, Γ IS CALLED THE BRONWICH CONTOUR.

PROOF LET $\hat{G}(k) = \int_{-\infty}^\infty g(t) e^{-ikt} dt.$ THEN $g(t) = \frac{1}{2\pi} \int_{-\infty}^\infty \hat{G}(k) e^{ikt} dk.$

NOW WE LET $g(t) = \begin{cases} e^{-\gamma t} f(t) & t > 0 \\ 0 & t < 0 \end{cases}$

WHERE $|f(t)| \leq C e^{\alpha t}$ AS $t \rightarrow \infty$, WHERE $\gamma > \alpha$. HENCE $\int_0^\infty g(t) dt < \infty$.

NOW, WE CALCULATE

$\hat{G}(k) = \int_0^\infty e^{-(\gamma+ik)t} f(t) dt.$ LET $s = \gamma + ik$. IT IS THEN

DEFINED FOR $RE(s) = \gamma > \alpha$. THE INVERSION FORMULA IS, FOR $t > 0$, THAT

$2\pi e^{-\gamma t} f(t) = \int_{-\infty}^\infty \hat{G}(k) e^{ikt} dk.$

AND THAT $\int_{-\infty}^\infty \hat{G}(k) e^{ikt} dk = 0$ FOR $t < 0$. THEN, WE LET $s = \gamma + ik, ds = idk$.

WE HAVE THAT

$f(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} e^{st} \hat{G}(s) ds$

$\hat{G}(s) = \int_0^\infty e^{-st} f(t) dt$

WHERE $\hat{G}(s)$ IS ANALYTIC FOR $RE(s) > \gamma$.

(ix) $\mathcal{L}(t^n) = \int_0^\infty e^{-st} t^n dt.$ LET $x = st \rightarrow dt = \frac{1}{s} dx \Rightarrow \mathcal{L}(t^n) = \frac{1}{s^{n+1}} \int_0^\infty e^{-x} x^n dx.$

NOW integration by parts gives $\int_0^\infty e^{-x} x^n dx = n!$.

$\rightarrow \mathcal{L}(t^n) = n! / s^{n+1}.$

TO CALCULATE $f(t)$ WE USE RESIDUE THEORY. SUPPOSE $F(s)$ HAS POLES IN $RE(s) < \gamma$: THEN, $f(t) = \frac{1}{2\pi i} \int_{\Gamma} F(s) e^{st} ds = -\lim_{R \rightarrow \infty} \int_{C_R} \frac{1}{2\pi i} F(s) e^{st} ds + \sum_{j=1}^N \text{RES}[F(s) e^{st}; s_j]$

THEREFORE, WHEN $F(s)$ HAS ONLY POLES IN REGION $RE(s) < \gamma$, WE CALCULATE

$$f(t) = \sum_{j=1}^N \text{RES}[F(s) e^{st}; s_j] - \lim_{R \rightarrow \infty} \int_{C_R} \frac{1}{2\pi i} F(s) e^{st} ds.$$

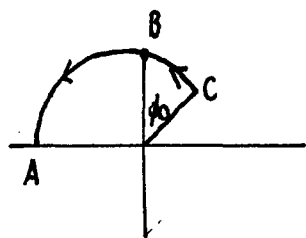
LEMMA SUPPOSE $\exists M > 0, K > 0$, SUCH THAT ON C_R ($s = Re^{i\phi}$),

$$|F(s)| \leq M/R^K.$$

THEN, WE HAVE

$$\lim_{R \rightarrow \infty} \left| \int_{C_R} e^{st} F(s) ds \right| = 0.$$

PROOF



ON CB: WE HAVE $s = Re^{i\phi}$ WITH $\phi_0 < \phi < \pi/2$.

$$\begin{aligned} \left| \int_{C_B} \right| &= \frac{1}{2\pi} \left| \int_{\phi_0}^{\pi/2} e^{Re^{i\phi}t} F(Re^{i\phi}) i Re^{i\phi} d\phi \right| \\ &\leq \frac{1}{2\pi} \int_{\phi_0}^{\pi/2} R e^{Rt \cos \phi} |F(Re^{i\phi})| d\phi \leq \frac{M}{2\pi R^{K-1}} \int_{\phi_0}^{\pi/2} e^{Rt \cos \phi} d\phi \end{aligned}$$

NOW LET $\phi = \pi/2 - \theta$. $\cos \theta = \cos(\pi/2 - \phi) = \sin \phi$. THEN, WITH $\phi_0 = \sin^{-1}(\gamma/R)$,

$$\left| \int_{C_B} \right| \leq \frac{M}{2\pi R^{K-1}} \int_0^{\phi_0} e^{Rt \sin \phi} d\phi \leq \frac{M}{2\pi R^{K-1}} \int_0^{\phi_0} e^{Rt \sin \phi_0} d\phi = \frac{M}{2\pi R^{K-1}} \int_0^{\phi_0} e^{\gamma t} dt$$

SINCE $\sin \phi \leq \sin \phi_0 = \gamma/R$. THEREFORE,

$$\left| \int_{C_B} \right| \leq \frac{M}{2\pi R^{K-1}} e^{\gamma t} \sin^{-1}\left(\frac{\gamma}{R}\right) \rightarrow \frac{M\gamma}{2\pi R^K} e^{\gamma t} \text{ AS } R \rightarrow \infty. \Rightarrow \left| \int_{C_B} \right| \rightarrow 0,$$

AS $R \rightarrow \infty$.

NOW ON AB WE CALCULATE

$$\left| \int_{C_{AB}} \right| \leq \frac{1}{2\pi} \int_{\pi/2}^{\pi} |e^{st}| |F(s)| |ds| \leq \frac{1}{2\pi} \int_{\pi/2}^{\pi} e^{Rt \cos \theta} |F| R d\theta$$

THEREFORE,

$$\left| \int_{C_{AB}} \right| \leq \frac{1}{2\pi} \frac{M}{R^{K-1}} \int_{\pi/2}^{\pi} e^{tR \cos \theta} d\theta = \frac{1}{2\pi} \frac{M}{R^{K-1}} \int_0^{\pi/2} e^{-Rt \sin \phi} d\phi. (\theta = \pi/2 + \phi).$$

THIS GIVES UPON USING $\sin \theta \geq 2\theta/\pi$ FOR $0 \leq \theta \leq \pi/2$, THAT

$$\left| \int_{C_{AB}} \right| \leq \frac{M}{2\pi R^{K-1}} \int_0^{\pi/2} e^{-2Rt \phi/\pi} d\phi = \frac{M}{R^{K-1} 2\pi} \left(\frac{\pi}{2Rt}\right) (1 - e^{-Rt}) \rightarrow 0 \text{ AS } R \rightarrow \infty, \text{ IF } t > 0.$$

EXAMPLE

SUPPOSE THAT $F(s) = \frac{P(s)}{Q(s)}$ WHERE $P(s), Q(s)$ POLYNOMIALS WITH $\deg Q > \deg P$.

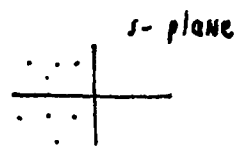
SUPPOSE THAT $Q(s_j) = 0$ $j=1, \dots, N$ AND $Q'(s_j) \neq 0$. HENCE s_j IS A SIMPLE ROOT.

NOW WE CALCULATE,

$$f(t) = Q^{-1}[F(s)] = \sum_{j=1}^N \text{RES} \left[\frac{P(s)e^{st}}{Q(s)}; s_j \right] = \sum_{j=1}^N \frac{P(s_j)}{Q'(s_j)} e^{s_j t}.$$

REMARKS

(i) IF ALL $\text{RE}(s_j) < 0 \Rightarrow f(t) \rightarrow 0$ AS $t \rightarrow \infty$



(ii) THE DOMINANT BEHAVIOR FOR LARGE t IS DETERMINED BY THE POLE WITH THE LARGEST REAL PART.

LET $\text{RE}(s_J) > \text{RE}(s_j)$ FOR $j=1, \dots, N, j \neq J$.

THEN, $f(t) \approx \frac{P(s_J)}{Q'(s_J)} e^{s_J t}$ AS $t \rightarrow \infty$.

EXAMPLE

(i) $F(s) = \frac{1}{s-a}, \Rightarrow f(t) = e^{at}$.

(ii) $F(s) = \frac{w}{s^2+w^2} \Rightarrow f(t) = \text{RES} \left[\frac{w}{s^2+w^2} e^{st}; iw \right] + \text{RES} \left[\frac{w}{s^2+w^2} e^{st}; -iw \right]$
 $f(t) = \frac{1}{2i} (e^{iwt} - e^{-iwt}) = \sin(\omega t)$.

(iii) $F(s) = \frac{s}{s^2+w^2} \Rightarrow f(t) = \cos(\omega t)$

(iv) $F(s) = \frac{1}{s(s^2+1)} = G(s)H(s), \quad g(t) = 1, \quad H(t) = \sin t. \quad f(t) = \int_0^t \sin(\tau) d\tau = -\cos \tau \Big|_0^t$
 THEN, $f(t) = 1 - \cos t$.

(v) $F(s) = \frac{1}{(s+1)(s-2)^2}$.

$$f(t) = \text{RES} \left[\frac{e^{st}}{(s+1)(s-2)^2}; -1 \right] + \text{RES} \left[\frac{e^{st}}{(s+1)(s-2)^2}; 2 \right]$$

$$= \frac{e^{-t}}{9} + \lim_{s \rightarrow 2} \left[\frac{d}{ds} \left[\frac{e^{st}}{(s+1)^2} F(s) \right] \Big|_{s=2} \right] = \lim_{s \rightarrow 2} \left[\frac{d}{ds} \left[(s+1)^{-1} e^{st} \right] \right] + e^{-t}/9.$$

CALCULATING WE GET $f(t) = e^{-t}/9 + \lim_{s \rightarrow 2} \left[-(s+1)^{-2} e^{st} + t(s+1)^{-1} e^{st} \right] = t e^{2t}/3$.

(vi) GIVEN $y'' - y' - 6y = 1 + e^{-t}$
 $y(0) = d, y'(0) = 0.$

FIND d SO THAT y IS BOUNDED AS $t \rightarrow \infty.$

TAKE LAPLACE TRANSFORM TO GET

$$(s^2 - s - 6)Y - s d + d = \frac{1}{s} + \frac{1}{s+1} \rightarrow Y(s) = \frac{d(s-1) + \frac{1}{s} + \frac{1}{s+1}}{(s-3)(s+2)}$$

WE OBTAIN, $Y(s) = \frac{d s (s^2 - 1) + 2s + 1}{s (s-3) (s+2) (s+1)} = \frac{P(s)}{Q(s)}$

WE HAVE POLES AT $s = 0, -1, -2, 3.$ THE POLE AT $s = 3$ GIVES $y \sim () e^{3t}$ AS $t \rightarrow \infty.$
 HENCE TO HAVE A BOUNDED SOLUTION AS $t \rightarrow \infty,$ WE REQUIRE

$$y \sim \frac{P(3)}{Q'(3)} e^{3t} + \frac{P(0)}{Q'(0)} e^{0t} + \dots \quad \text{i.e. } P(3) = 0.$$

THIS GIVES, $24d + 7 = 0 \rightarrow d = -7/24.$

(vii) GIVEN $F(s) = \frac{s}{(s^2+1)^2},$ FIND $f(t).$ WE NOTICE THAT

$$F(s) = -\frac{d}{ds} G(s), \quad \text{WHERE } G(s) = \frac{1}{2} (s^2+1)^{-1} \rightarrow g(t) = \frac{1}{2} \sin t.$$

RECALLING $d[t g(t)] = -G'(s) \rightarrow d^{-1}[F(s)] = +\frac{t}{2} \sin t.$

(viii) LET $F(s) = \frac{1}{(s^2+1)^2}$ pole of order 2 at $s = \pm i.$

$$\text{RES} [; i] = \lim_{s \rightarrow i} \left[\frac{d}{ds} \left(\frac{e^{st}}{(s+i)^2} \right) \right] = \lim_{s \rightarrow i} \left[\frac{t e^{st}}{(s+i)^2} - \frac{2 e^{st}}{(s+i)^3} \right] = \frac{-t e^{it}}{4} + \frac{e^{it}}{4i}$$

$$\text{RES} [; -i] = \lim_{s \rightarrow -i} \left[\frac{d}{ds} \left(\frac{e^{st}}{(s-i)^2} \right) \right] = \lim_{s \rightarrow -i} \left[\frac{t e^{st}}{(s-i)^2} - \frac{2 e^{st}}{(s-i)^3} \right] = \frac{-t e^{-it}}{4} - \frac{1}{4i} e^{-it}$$

THIS GIVES, $f(t) = -\frac{t}{4} (e^{it} + e^{-it}) + \frac{1}{4i} (e^{it} - e^{-it})$

$$f(t) = -\frac{t}{2} \cos t + \frac{1}{2} \sin t.$$

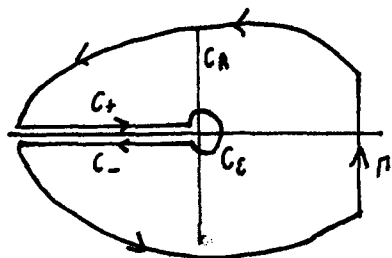
(ix) LET $X(s) = \frac{1}{s^2(s^2+w^2)}.$ $F(s) = 1/s^2 \rightarrow f(t) = t$

$$G(s) = \frac{1}{s^2+w^2} \rightarrow g(t) = \frac{1}{w} \sin(wt)$$

THEN $h(t) = d^{-1}[X(s)] = \int_0^t \frac{t}{w} \sin(w(t-\tau)) d\tau = \frac{t}{w^2} - \frac{\sin(wt)}{w^3}$ integration by parts.

EXAMPLE 1

WE LET $F(s) = s^{-a}$ WITH $0 < a < 1$. WE REQUIRE $F(s)$ ANALYTIC FOR $\text{Re}(s) > \chi$. HENCE, WE CAN TAKE THE CONTOUR AS SHOWN



SO $f(t) = \frac{1}{2\pi i} \int_{\Gamma} s^{-a} e^{st} ds$.

NOW ON C_E : $|\frac{1}{2\pi i} \int_{C_E} s^{-a} e^{st} ds| \leq C \epsilon^{1-a} \rightarrow 0$ AS $\epsilon \rightarrow 0$.

NOW $|\int_{C_R}| \rightarrow 0$ AS $R \rightarrow \infty$, SINCE $|F| \rightarrow 0$ AS $R \rightarrow \infty$.

NOW WE HAVE, $\frac{1}{2\pi i} \int_{\Gamma} + \frac{1}{2\pi i} (\int_{C+} + \int_{C-}) = 0$.

THEREFORE, $f(t) = -\frac{1}{2\pi i} \int_{C+} - \frac{1}{2\pi i} \int_{C-}$.

NOW ON $C+$: $s = re^{i\pi}$ so $ds = -dr$. $\rightarrow \frac{1}{2\pi i} \int_{C+} = \frac{1}{2\pi i} \int_{\infty}^0 e^{rt} r^{-a} e^{-i\pi a} (-dr)$

NOW ON $C-$: $s = re^{-i\pi}$ so $ds = -dr$. $\rightarrow \frac{1}{2\pi i} \int_{C-} = \frac{1}{2\pi i} \int_0^{\infty} e^{-rt} r^{-a} e^{i\pi a} (-dr)$.

THEREFORE, $f(t) = -\frac{1}{2\pi i} \int_0^{\infty} r^{-a} (e^{-i\pi a} - e^{i\pi a}) e^{-rt} dr = \frac{1}{\pi} \int_0^{\infty} r^{-a} \sin(\pi a) e^{-rt} dr$

NOW $f(t) = \frac{\sin(\pi a)}{\pi} \int_0^{\infty} r^{-a} e^{-rt} dr$. LET $x = rt$ $dr = \frac{1}{t} dx$
 $r^{-a} = x^{-a} t^a$.

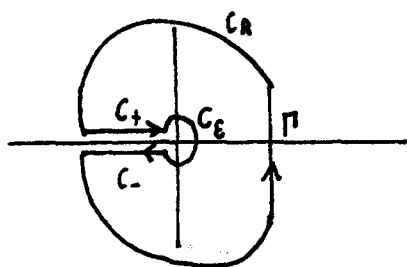
HENCE, $f(t) = t^{a-1} \frac{\sin(\pi a)}{\pi} \int_0^{\infty} x^{-a} e^{-x} dx$.

RECALL $\Gamma(z) = \int_0^{\infty} x^{z-1} e^{-x} dx$ is the gamma function. Therefore,

$f(t) = t^{a-1} \frac{\sin(\pi a)}{\pi} \Gamma(1-a)$.

EXAMPLE 2 FIND THE INVERSE OF $F(s) = \frac{e^{-a\sqrt{s}}}{s}$ WHERE $a \geq 0$.

WE TAKE THE CONTOUR AS SHOWN:



WE HAVE $f(t) = \frac{1}{2\pi i} \int_{\Gamma} \frac{e^{-a\sqrt{s} + st}}{s} ds$.

NOW $f(\frac{1}{t}) = -\frac{1}{2\pi i} \int_{C+} - \frac{1}{2\pi i} \int_{C-} - \lim_{\epsilon \rightarrow 0} \frac{1}{2\pi i} \int_{C_E}$.

NOW WE CALCULATE

7

$$\lim_{\epsilon \rightarrow 0} \left[\frac{-1}{2\pi i} \oint_{C_\epsilon} \frac{e^{-a\sqrt{s}}}{s} e^{st} ds \right] \rightarrow \frac{-1}{2\pi i} \oint_{C_\epsilon} \frac{1}{s} ds = 1.$$

NOW ON C_+ : $s = re^{i\pi}$ $s^{1/2} = r^{1/2} e^{i\pi/2} = i r^{1/2}$.

C_- : $s = re^{-i\pi}$ $s^{1/2} = r^{1/2} e^{-i\pi/2} = -i r^{1/2}$.

THUS
$$f(t) = 1 - \frac{1}{2\pi i} \int_{\infty}^0 \frac{e^{-rt} e^{-iar^{1/2}}}{-r} (-dr) - \frac{1}{2\pi i} \int_0^{\infty} \frac{e^{-rt} e^{iar^{1/2}}}{-r} (-dr).$$

HENCE,
$$f(t) = 1 - \frac{1}{2\pi i} \int_0^{\infty} \frac{e^{-rt}}{r} (e^{iar^{1/2}} - e^{-iar^{1/2}}) dr$$

WE CALCULATE,

$$f(t) = 1 - \frac{1}{\pi} \int_0^{\infty} \frac{e^{-rt}}{r} \sin(ar^{1/2}) dr$$

LET $u = r^{1/2}$, $r = u^2$, $dr = 2u du$. THUS

$$f(t) = 1 - \frac{2}{\pi} \int_0^{\infty} e^{-u^2 t} \frac{\sin(au)}{u} du.$$

WE DEFINE
$$I(a) = \frac{2}{\pi} \int_0^{\infty} e^{-u^2 t} \frac{\sin(au)}{u} du.$$

WE GET,

$$I'(a) = \frac{2}{\pi} \int_0^{\infty} e^{-u^2 t} \cos(au) du = \frac{2}{\pi} \operatorname{RE} \left[\int_0^{\infty} e^{-u^2 t + iau} du \right]$$

$$I'(a) = \operatorname{RE} \left[\frac{1}{\sqrt{t}} \int_{-\infty}^{\infty} e^{-t(u^2 - iau/t - a^2/4t^2) - a^2/4t} du \right] = \frac{e^{-a^2/4t}}{\pi} \operatorname{RE} \left[\int_{-\infty}^{\infty} e^{-t(u - i a/2t)^2} du \right]$$

NOW LET $2\sigma^2 = 1/t$, SO THAT $\sigma = (1/2t)^{1/2}$. HENCE,

$$I'(a) = \frac{1}{\pi} e^{-a^2/4t} \sqrt{2\pi} \sigma = \frac{1}{\sqrt{\pi t}} e^{-a^2/4t}.$$

WE CALCULATE, $I(a) = \frac{1}{\sqrt{\pi t}} \int_0^a e^{-\lambda^2/4t} d\lambda$. LET $\zeta = \lambda/2t^{1/2}$. $d\lambda = 2t^{1/2} d\zeta$.

THUS
$$I(a) = \frac{2}{\sqrt{\pi}} \int_0^{a/2t^{1/2}} e^{-\zeta^2} d\zeta = \operatorname{ERF} \left(a/2t^{1/2} \right), \text{ WHERE } \operatorname{ERF}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-\zeta^2} d\zeta.$$

FINALLY,
$$f(t) = 1 - \operatorname{ERF} \left(a/2\sqrt{t} \right).$$

$$1. \quad \text{ERF}(z) \equiv \frac{2}{\sqrt{\pi}} \int_0^z e^{-\zeta^2} d\zeta. \quad \text{NOW } \text{ERF}(0) = 0, \quad \text{ERF}(\infty) = 1$$

$$\text{NOW } \text{ERFC}(z) = \frac{2}{\sqrt{\pi}} \int_z^\infty e^{-\zeta^2} d\zeta, \quad \text{ERFC}(\infty) = 0, \quad \text{ERFC}(0) = 1.$$

$$\text{ALSO } \text{ERF}(z) + \text{ERFC}(z) = 1.$$

$$2. \quad \text{THEREFORE IF } F(s) = \frac{e^{-a\sqrt{s}}}{s}$$

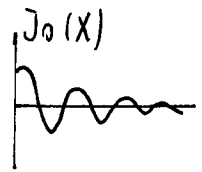
$$f(t) = 1 - \text{ERF}\left(\frac{a}{2\sqrt{t}}\right) = \text{ERFC}\left(\frac{a}{2\sqrt{t}}\right).$$

PROBLEM

CONSIDER

$$xy'' + y' + xy = 0 \quad x \geq 0$$

$$y(0) = 1, \quad y'(0) \text{ FINITE}$$



THE SOLUTION IS $y = J_0(x)$ WHERE $J_0(x)$ IS THE BESSEL FUNCTION OF ORDER ZERO OF FIRST KIND. WE WILL FIND AN INTEGRAL REPRESENTATION OF $J_0(x)$ USING A LAPLACE TRANSFORM SOLUTION.

$$\mathcal{L}(xy'') + \mathcal{L}(y') + \mathcal{L}(xy) = 0 \quad \text{BUT } \mathcal{L}(xy'') = \int_0^\infty xy'' e^{-sx} dx$$

$$\rightarrow -\frac{d}{ds} [s^2 Y - sy(0) - y'(0)] + sY - y(0) - Y' = 0$$

$$\text{SO } \mathcal{L}(xy'') = -\frac{d}{ds} \int_0^\infty y'' e^{-sx} dx$$

$$= -\frac{d}{ds} [s^2 Y(s) - sy(0) - y'(0)]$$

THIS YIELDS THAT

$$\frac{d}{ds} (s^2 Y) + Y' - sY = 0$$

$$\text{ALSO } \mathcal{L}(xy) = \int_0^\infty xye^{-sx} dx = -\frac{d}{ds} \int_0^\infty ye^{-sx} dx$$

$$\text{OR EQUIVALENTLY } (s^2 + 1)Y' + sY = 0 \rightarrow \frac{Y'}{Y} = -\frac{s}{s^2 + 1} \rightarrow (\ln Y)' = -\frac{1}{2} \log(s^2 + 1)$$

HENCE $Y(s) = \frac{C}{\sqrt{s^2 + 1}}$, where C IS ARBITRARY AT THIS STAGE.

$$\text{SO } \frac{C}{(s^2 + 1)^{1/2}} = \int_0^\infty y(x) e^{-sx} dx$$

NOW $y(0) = 1$ $x \rightarrow 0$ CORRESPONDS TO $s \rightarrow \infty$.

$$\text{HENCE FOR } s \gg 1 \quad \frac{C}{s} \sim \int_0^\infty e^{-sx} dx = \frac{1}{s}.$$

$$\text{THIS GIVES } C = 1, \text{ WHICH IMPLIES } J_0(x) = \frac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} \frac{e^{sx}}{\sqrt{s^2 + 1}} ds. \quad (a > 0)$$

PROBLEM 1 CONSIDER $U_t = D U_{xx} \quad 0 < x < \infty, t > 0$

$$U(x,0) = 0, \quad U(0,t) = U_0, \quad U_0 \text{ CONSTANT.}$$

NOW TAKE LAPLACE TRANSFORMS $\mathcal{U}(x,s) = \mathcal{L}(U)$. NOW WE CALCULATE,

$$\mathcal{L}(U_t) = D \mathcal{L}(U_{xx}), \quad \text{WITH } \mathcal{U}(x,s) = \mathcal{L}(U(x,t)).$$

WE CALCULATE,

$$s \mathcal{U} = D \mathcal{U}_{xx} \quad \rightarrow \quad \mathcal{U} = a_0 e^{-\sqrt{s}x/\sqrt{D}} + b_0 e^{\sqrt{s}x/\sqrt{D}}$$

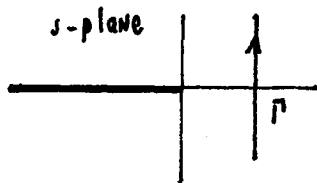
$$\mathcal{U}(0,s) = U_0/s$$

WE CALCULATE

$$\mathcal{U}(x,s) = \frac{U_0}{s} e^{-\sqrt{s/\sqrt{D}} x}$$

NOW WE NEED $\text{RE}(\sqrt{s}) > 0$. $\sqrt{s} = \sqrt{r} e^{i\theta/2} \rightarrow \omega(\theta/2) > 0 \rightarrow -\pi < \theta < \pi$.

THU, WE TAKE BRANCH CUT ON NEGATIVE REAL AXIS.



$$U(x,t) = \frac{U_0}{2\pi i} \int_{\Gamma} \frac{1}{s} e^{-\sqrt{s/\sqrt{D}} x} e^{st} ds$$

WE CALCULATED THU EARLIER, WITH THE RESULT

$$U(x,t) = U_0 \text{ERFC} \left(\frac{x}{2\sqrt{Dt}} \right).$$

PROBLEM 2 WE WANT TO SOLVE

$$U_t = D U_{xx} \quad -\infty < x < \infty, t > 0$$

$$U(x,0) = \delta(x - x_0)$$

WE CALCULATE THE LAPLACE TRANSFORM, $\mathcal{U}(x,s) = \int_0^\infty e^{-st} U(x,t) dt$

$$s \mathcal{U} - \delta(x - x_0) = D \mathcal{U}_{xx}$$

WE HAVE

$$D \mathcal{U}_{xx} - s \mathcal{U} = -\delta(x - x_0) \quad -\infty < x < \infty$$

\mathcal{U} BOUNDED AS $|x| \rightarrow \infty$.

THE CONTINUITY AND JUMP CONDITIONS ARE $[\mathcal{U}]|_{x_0} = 0, \quad [D \mathcal{U}_x]|_{x_0} = -1$.

NOW WE GET

$$\mathcal{U}(x,s) = \begin{cases} A e^{+\sqrt{s/D} (x-x_0)} & x < x_0 \\ A e^{-\sqrt{s/D} (x-x_0)} & x > x_0. \end{cases}$$

$$\text{NOW} \quad D \mathcal{U}_x|_{x_0^+} - D \mathcal{U}_x|_{x_0^-} = -1 \quad \rightarrow \quad -2 \sqrt{\frac{s}{D}} A = -\frac{1}{D} \quad \rightarrow \quad A = \frac{1}{2\sqrt{Ds}}$$

NOW WE CALCULATE,

$$U(x,s) = \frac{1}{2\sqrt{Ds}} e^{-\sqrt{s/D} |x-x_0|}$$

NOW RECALL, $\mathcal{L}^{-1} \left[\frac{e^{-a\sqrt{s}}}{\sqrt{s}} \right] = \frac{1}{\sqrt{\pi t}} e^{-a^2/4t}$, WHERE $a = |x-x_0|/\sqrt{D}$.

HENCE, $U(x,t) = \frac{1}{2\sqrt{\pi Dt}} e^{-(x-x_0)^2/4Dt}$.

EXAMPLE 3 NOW CONSIDER

$$U_t = D U_{xx} \quad 0 < x < 1$$

$$U_x(0,t) = 0, \quad U(1,t) = 1, \quad U(x,0) = 0.$$

NOW WE TAKE LAPLACE TRANSFORMS TO GET

$$D U_{xx} - s U = 0 \quad 0 < x < 1$$

$$U_x = 0 \text{ ON } x=0, \quad U = \frac{1}{s} \text{ ON } x=1.$$

WE CALCULATE,

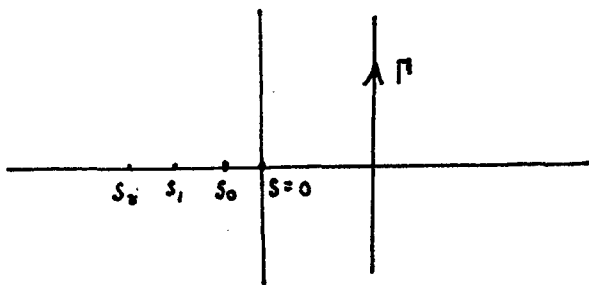
$$U(x,s) = \frac{1}{s} \frac{\cosh(\sqrt{s/D} x)}{\cosh(\sqrt{s/D})} \quad U \text{ has only poles, no branch cuts.}$$

NOW $s=0$ IS A SIMPLE POLE. ALSO POLES OF $\cosh(\sqrt{s/D}) = 0$. HENCE

$$\sqrt{s/D} = (k + \frac{1}{2})\pi i, \text{ THEN } s_k = -D(k + \frac{1}{2})^2 \pi^2 \quad k=0,1,2,\dots$$

NOTICE $s_0 = s_{-1}, \quad s_1 = s_{-2}$.

NOW WE USE RESIDUE THEOREM



$$U(x,t) = \frac{1}{2\pi i} \int_{\Gamma} \frac{1}{s} \frac{\cosh(\sqrt{s/D} x)}{\cosh(\sqrt{s/D})} e^{st} ds$$

$$U(x,t) = \text{RES} [; s=0] + \sum_{k=0}^{\infty} \text{RES} \left[\frac{1}{s} \frac{\cosh(\sqrt{s/D} x)}{\cosh(\sqrt{s/D})} e^{st} ; s_k \right]$$

$$U(x,t) = 1 + \sum_{k=0}^{\infty} \frac{e^{s_k t} \cosh[(k + \frac{1}{2})\pi i x]}{s_k \frac{1}{2} \frac{1}{(s_k D)^{1/2}} \sinh[(k + \frac{1}{2})\pi i]} = 1 + 2 \sum_{k=0}^{\infty} \frac{\cos((k + \frac{1}{2})\pi x) e^{s_k t}}{(k + \frac{1}{2})\pi i \cdot i \cdot (-1)^k}$$

THEREFORE,

$$U = 1 + \frac{2}{\pi} \sum_{k=0}^{\infty} \frac{(-1)^{k+1}}{(k + \frac{1}{2})} \cos((k + \frac{1}{2})\pi x) e^{-D(k + \frac{1}{2})^2 \pi^2 t}$$

NOW FOR $t \gg 1$ WE GET $U(x,t) \sim 1 - \frac{4}{\pi} e^{-D^2 \pi^2 t / 4} \cos(\pi x / 2)$.

NOW LARGE s CORRESPONDS TO SMALL t . FOR $s \gg 1$ WE USE $\cosh(z) \sim e^z / 2$, WE OBTAIN

$$U(x,s) \sim \frac{1}{s} e^{-\sqrt{s/D}(1-x)}$$

LETTING $a = (1-x)/\sqrt{D}$ AND RECALLING $\mathcal{L}^{-1} \left[\frac{1}{s} e^{-a\sqrt{s}} \right] = \text{erfc} \left(\frac{a}{2\sqrt{t}} \right)$.

THUS, $U(x,t) \sim \text{erfc} \left(\frac{(1-x)}{2\sqrt{Dt}} \right)$ FOR $t \ll 1$.

EXAMPLE (WAVE EQUATION)

CONSIDER $U_{tt} = U_{xx} \quad 0 < x < L, t > 0$
 $U(x,0) = U_t(x,0) = 0$
 $0 = U(0,t); U(L,t) = 1.$

WE HAVE TWO REPRESENTATIONS OF THE SOLUTION:

METHOD 1 WE TAKE LAPLACE TRANSFORMS TO OBTAIN,

$$U_{xx} - s^2 U = 0 \quad 0 < x < L$$
$$U(0,s) = 0, \quad U(L,s) = \frac{1}{s}.$$

WE SOLVE TO GET $U(x,s) = a_0 \cosh(sx) + b_0 \sinh(sx)$. SATISFYING THE BOUNDARY CONDITION WE OBTAIN

$$U(x,s) = \frac{1}{s} \frac{\sinh(sx)}{\sinh(sL)}$$

THE POLES ARE WHERE $sL = n\pi i$ WITH $n = \pm 1, \pm 2, \dots$ ALSO $n = 0$.

THE REFERENCE,

$$U(x,t) = \text{RES} \left[\frac{1}{s} \frac{\sinh(sx)}{\sinh(sL)} e^{st}; 0 \right] + \sum_{n=1}^{\infty} \text{RES} \left[\frac{1}{s} \frac{\sinh(sx)}{\sinh(sL)} e^{st}; \frac{n\pi i}{L} \right] + \sum_{n=1}^{\infty} \text{RES} \left[\frac{1}{s} \frac{\sinh(sx)}{\sinh(sL)} e^{st}; \frac{-n\pi i}{L} \right]$$

WE NOW CALCULATE

$$U(x,t) = x/L + \sum_{n=1}^{\infty} \frac{L}{n\pi i} \frac{\sinh(n\pi i x/L)}{L \cosh(n\pi i)} e^{n\pi i t/L} + \sum_{n=1}^{\infty} \frac{L}{n\pi i} \frac{\sinh(n\pi i x/L)}{L \cosh(n\pi i)} e^{-n\pi i t/L}$$

NOW $\cos(n\pi i) = (-1)^n$. WE OBTAIN,

$$u(x,t) = x/L + \sum_{n=1}^{\infty} \frac{(-1)^n}{n\pi} \sin\left(\frac{n\pi x}{L}\right) e^{n\pi i t/L} + \sum_{n=1}^{\infty} \frac{(-1)^n}{n\pi} \sin\left(\frac{n\pi x}{L}\right) e^{-n\pi i t/L}$$

↑ replace n by $-n'$.

THEN, WE CALCULATE

$$u(x,t) = x/L + 2 \sum_{n=1}^{\infty} \frac{(-1)^n}{n\pi} \sin\left(\frac{n\pi x}{L}\right) \cos\left(\frac{n\pi t}{L}\right)$$

METHOD 2

$$u(x,s) = \frac{1}{s} \frac{\sinh(sx)}{\sinh(sL)} = \frac{1}{s} \left[\frac{e^{sx} - e^{-sx}}{e^{sL} - e^{-sL}} \right] = \frac{1}{s} \frac{e^{sx}}{e^{sL}} \left(\frac{1 - e^{-2sx}}{1 - e^{-2sL}} \right)$$

NOW USING $\frac{1}{1-x} = 1 + x + x^2 + x^3 + \dots$ FOR $|x| < 1$, WE CALCULATE

$$u(x,s) = \frac{1}{s} e^{-s(L-x)} [1 - e^{-2sx}] [1 + e^{-2sL} + e^{-4sL} + \dots + e^{-2nsL} + \dots]$$

THIS GIVES,

$$u(x,s) = \frac{1}{s} \left[e^{-s(L-x)} - e^{-s(L+x)} \right] [1 + e^{-2sL} + e^{-4sL} + \dots + e^{-2nsL} + \dots]$$

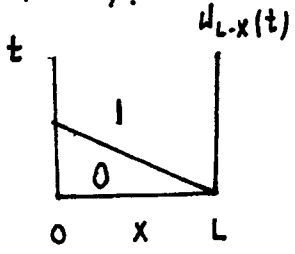
$$u(x,s) = \frac{1}{s} \left[e^{-s(L-x)} + e^{-s(3L-x)} + \dots + e^{-s((2n+1)L-x)} + \dots \right] - \frac{1}{s} \left[e^{-s(L+x)} + e^{-s(3L+x)} + \dots + e^{-s((2n+1)L+x)} \right]$$

NOW RECALL $\mathcal{L}(u_a(t)) = \frac{1}{s} e^{-as}$ WITH $u_a(t) = \begin{cases} 1 & \text{if } t > a \\ 0 & \text{if } 0 < t < a \end{cases}$

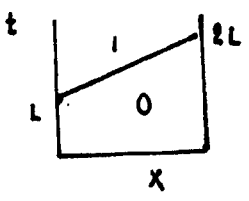
HENCE, $u(x,t) = (u_{L-x}(t) + u_{3L-x}(t) + \dots + u_{(2n+1)L-x}(t) + \dots) - (u_{L+x}(t) + u_{3L+x}(t) + \dots + u_{(2n+1)L+x}(t) + \dots)$
 (Reflected wave)

REMARK

$$u_{L-x}(t) = \begin{cases} 1 & \text{if } t > L-x \rightarrow x > L-t \\ 0 & \text{if } 0 < t < L-x \end{cases}$$

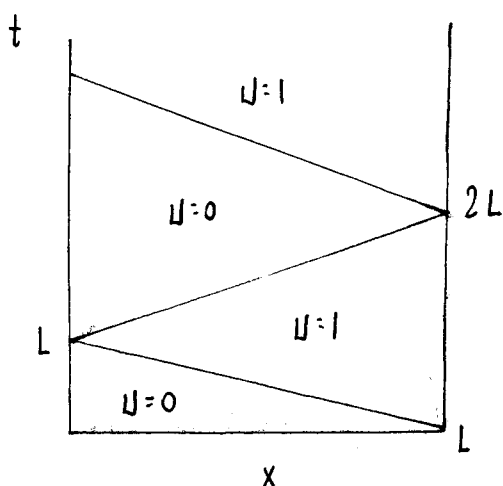


$$u_{L+x}(t) = \begin{cases} 1 & \text{if } t > L+x \\ 0 & \text{if } t < L+x \end{cases}$$



THEREFORE THE SOLUTION HAS THE FORM $u = u_{L-x}(t) - u_{L+x}(t) + \dots$

(13)



EXAMPLE FIND AN INTEGRAL REPRESENTATION FOR THE SOLUTION TO THE WAVE

EQUATION $u_{tt} = u_{xx}$

$$u(x,0) = u_t(x,0) = 0, \quad u(0,t) = 0, \quad u(L,t) = f(t)$$

THIS CORRESPONDS TO BOUNDARY FORCING WHICH GENERATES WAVES IN $0 < x < L$.

SOLUTION WE TAKE $\bar{u}(x,s) = \int_0^\infty e^{-st} u(x,t) dt$. NOW WITH $u = u_t = 0$ AT $t=0$

THIS GIVES

$$s^2 \bar{u} = \bar{u}_{xx}$$

$$\bar{u}(0,s) = 0, \quad \bar{u}(L,s) = F(s) \quad \text{WHERE} \quad F(s) = \mathcal{L}\{f(t)\}$$

THE SOLUTION IS $\bar{u}(x,s) = A \sinh(sx)$ SO $F(s) = A \sinh(sL)$,

WHICH YIELDS

$$\bar{u}(x,s) = F(s) \frac{\sinh(sx)}{\sinh(sL)}$$

WE WRITE THIS AS $\bar{u}(x,s) = F(s) G(s)$ $G(s) = s \Phi(x,s)$

WHERE $\Phi(x,s) = \frac{1}{s} \frac{\sinh(sx)}{\sinh(sL)}$, THE INVERSE TRANSFORM $\phi(x,t) = \mathcal{L}^{-1}\{\Phi(x,s)\}$

WAS FOUND IN PREVIOUS PROBLEM. HENCE $\mathcal{L}^{-1}[s \Phi(x,s)] = \frac{d}{dt} \phi(x,t)$ SINCE $\phi(x,0) = 0$

THEN USING CONVOLUTION THEOREM WITH $\mathcal{L}^{-1}[F(s)] = f(t)$, $\mathcal{L}^{-1}[G(s)] = \frac{d}{dt} \phi(x,t)$

WE OBTAIN
$$u(x,t) = \int_0^t \left(\frac{d}{d\tau} \phi(x,\tau) \right) f(t-\tau) d\tau.$$

DIFFUSION IN SEMI-INFINITE MEDIUM

CONSIDER $U_t = D U_{xx} \quad 0 < x < \infty, t > 0$

$U(0, t) = h(t), \quad U(x, 0) = 0, \quad U \rightarrow 0 \text{ as } x \rightarrow \infty \text{ FIXED } t.$

NOW TAKE LAPLACE TRANSFORMS

$sU = D U_{xx}$ so $U(x, s) = A e^{-\sqrt{s/D} x} + B e^{\sqrt{s/D} x}$

$U(0, s) = H(s)$

WE SET $B = 0$ SO THAT $U \rightarrow 0$ AS $x \rightarrow \infty$.

THEN $U(0, s) = H(s)$ GIVES

$U(x, s) = H(s) e^{-\sqrt{s/D} x}$

NOW $\mathcal{L}^{-1}(e^{-a\sqrt{s}}) = \frac{a}{2\sqrt{\pi} t^{3/2}} e^{-a^2/4t}$

SO THAT $\mathcal{L}^{-1}(e^{-\sqrt{s/D} x}) = \frac{x}{2\sqrt{\pi D} t^{3/2}} e^{-x^2/4Dt}$

$\mathcal{L}^{-1}(H(s)) = h(t)$

AND THE CONVOLUTION THEOREM GIVES

$U(x, t) = \int_0^t h(t-\tau) \frac{x}{2\sqrt{\pi D} \tau^{3/2}} e^{-x^2/4D\tau} d\tau$

DELAY-DIFFERENTIAL EQUATIONS

CONSIDER $y'(t) = -a y(t-1) \quad 0 \leq t < \infty$

WITH $y(t) = 1$ FOR $-1 \leq t \leq 0$.

THIS IS A DELAY EQUATION. WHAT IS THE BEHAVIOR AS $t \rightarrow \infty$ FOR DIFFERENT VALUES OF a .

THEN $Y(s) = \mathcal{L}(y(t))$

$$sY(s) - y(0) = -a \int_0^\infty y(t-1)e^{-st} dt = -a \int_0^1 y(t-1)e^{-st} dt - a \int_1^\infty y(t-1)e^{-st} dt$$

so $sY(s) - 1 = -a \int_{-1}^0 y(s)e^{-s(s+1)} ds - a \left(\int_0^\infty y(s)e^{-ss} \right) e^{-s}$

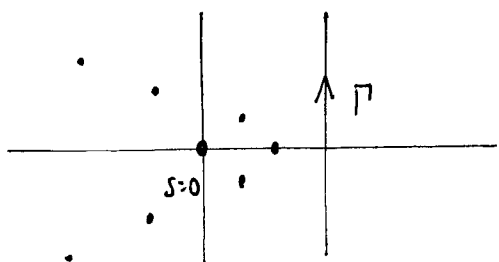
THIS GIVES, $sY(s) - 1 = -a \left(\int_{-1}^0 e^{-ss} ds \right) e^{-s} - a e^{-s} Y(s)$
 $= -a e^{-s} \left(-\frac{1}{s} e^{-ss} \Big|_{-1}^0 \right) - a e^{-s} Y(s)$

so $(s + a e^{-s}) Y(s) = 1 + \frac{a}{s} e^{-s} (1 - e^{s}) = 1 + \frac{a}{s} e^{-s} - \frac{a}{s}$

THIS GIVES $Y(s) = \frac{s - a + a e^{-s}}{s(s + a e^{-s})}$

NOTICE THAT THE INVERSE TRANSFORM IS $y(t) = \frac{1}{2\pi i} \int_{\Gamma} Y(s) e^{st} ds$

WITH ALL ZEROS OF $s(s + a e^{-s}) = 0$ TO THE LEFT OF VERTICAL LINE Γ .



THERE IS A POLE AT $s=0$ AND AT THE ZEROS OF $p(s) = s + a e^{-s} = 0$.

REMARK IF WE INITIALLY SUBSTITUTED $y = e^{\lambda t}$ WE WOULD OBTAIN $\lambda = -a e^{-\lambda}$

SO THAT λ IS A ROOT OF $p(\lambda) = \lambda + a e^{-\lambda} = 0$ FOR STABILITY WE NEED $\text{Re } \lambda < 0$.

EXAMPLE SUPPOSE $a = -1$. FIND BEHAVIOR AS $t \rightarrow \infty$.

WHEN $a = -1$ WE HAVE $p(s) = s - e^{-s} = 0$. WE NEED THE POLE WITH THE LARGEST REAL PART.

WE LET $s = x + iy$ SO THAT $x + iy = e^{-x} (\cos y - i \sin y)$

SO $x = e^{-x} \cos y$ $y = -e^{-x} \sin y$

NOW THE POLE WITH LARGEST REAL PART x HAS $x e^{+x}$ MAXIMIZED. THIS MEANS THAT $y = 0$ AND $x e^x = 1 \Rightarrow x = .567$ BY NEWTON'S METHOD.

HENCE LET $s_x = x_x$ SO THAT

$$y(t) = \frac{1}{2\pi i} \int_{\Gamma} Y(s) e^{st} ds \sim \text{RES} \left[\frac{s+1-e^{-s}}{s(s+ae^{-s})}; s_x \right] e^{s_x t} \quad \text{AS } t \rightarrow \infty.$$

HENCE
$$y(t) \sim \left(\frac{s_x + 1 - e^{-s_x}}{s_x (1 + e^{-s_x})} \right) e^{s_x t} = \frac{e^{s_x t}}{s_x (1 + e^{-s_x})} \quad \text{SINCE } s_x e^{-s_x} = 1$$

SO
$$y(t) \sim \frac{e^{s_x t}}{s_x (1 + e^{-s_x})} \quad \text{AS } t \rightarrow \infty \quad \text{WHERE } s_x = .567$$

(THIS GIVES $y(t) \sim 1.125 e^{.567 t}$)
AS $t \rightarrow \infty$

EXAMPLE (OSCILLATIONS)

NOW SUPPOSE $a > 0$. WHAT CAN WE SAY ABOUT THE ZEROS OF

$$p(s) = s + a e^{-s} = 0.$$

UNDER WHAT CONDITIONS ARE $\text{RE}(s) \leq 0$?

SUPPOSE WE LOOK FOR s ON IMAGINARY AXIS. WE SET $s = iy$ AND LET $y > 0$ (WLOG SINCE $s = -iy$ WOULD ALSO WORK).

THEN
$$p(iy) = iy + a e^{iy} = iy + a (\cos y + i \sin y) = 0$$

$$p(iy) = a \cos y + i(y - a \sin y)$$

NOTICE $\text{RE } p = 0$ WHEN $y = \pi/2, 3\pi/2, \dots$

$\text{IM } p = 0$ WHEN $y - a \sin y = 0 \rightarrow a = \frac{\pi}{2}$ WHEN $y = \frac{\pi}{2}$.

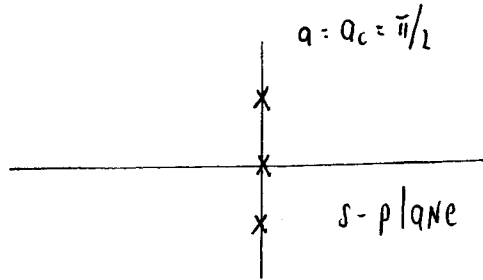
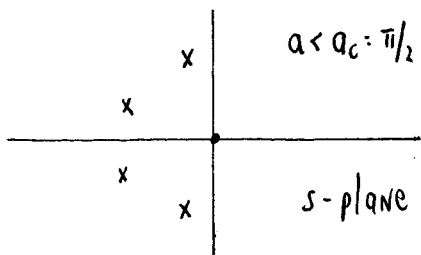
THEREFORE, WHEN $a = \pi/2$ WE HAVE $s = \pm i\pi/2$ ROOTS ON IMAGINARY AXIS.

WE WOULD LIKE TO PROVE THAT IF $0 < a < a_c = \pi/2$ WE HAVE $\text{RE } s < 0$

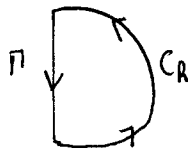
WHEREAS FOR $a = a_c = \pi/2 \rightarrow s = \pm i\pi/2$. HENCE WE CAN GET OSCILLATION FROM

$$y'(t) = -\frac{\pi}{2} y(t-1)$$

AS $t \rightarrow \infty$.



NOW TO PROVE THIS WE USE THE WINDING NUMBER CRITERION



$$N_0(p) = \frac{1}{2\pi} \left[\Delta_{\Gamma} \arg(p(iy)) + \lim_{R \rightarrow \infty} \Delta_{C_R} \arg(p(s)) \right]$$

$N_0(p) = \# \text{ ZEROS OF } p \text{ IN } \text{RE}(s) > 0.$

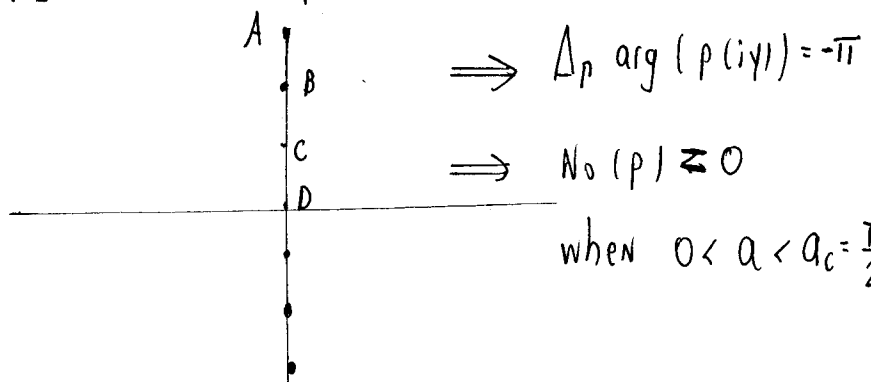
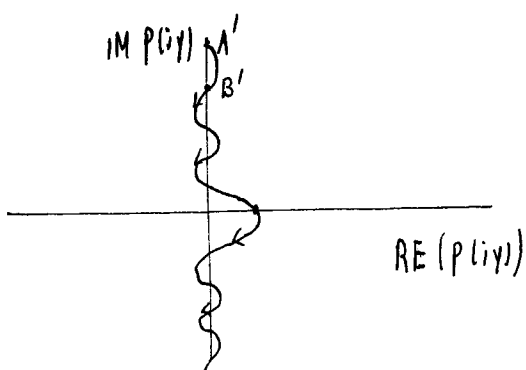
NOW $p(s) = s \left(1 + \frac{ae^{-s}}{s}\right)$ WITH $s = Re^{i\theta}$ SO $\lim_{R \rightarrow \infty} \Delta_{C_R} \arg(p(s)) = \pi$.

HENCE
$$N_0(p) = \frac{1}{2\pi} \left[\pi + \Delta_{\Gamma} \arg(p(s)) \right]$$

NOW ON Γ , $s = iy$ SO $p(iy) = a \cos y + i(y - a \sin y)$.

SUPPOSE $0 < a < a_c = \pi/2$. THEN IF $\cos y = 0$ WE HAVE $\text{IM}[p(iy)] = y - a \sin y > 0$

THUS $\text{RE}[p(iy)] = 0 \rightarrow \text{IM}[p(iy)] > 0$ FOR $y > 0$.



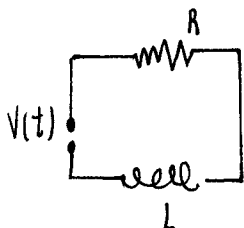
when $0 < a < a_c = \frac{\pi}{2}$

SUPPOSE THAT f IS PERIODIC WITH $f(t) = f(t+T)$.

THEN WE SHOWED EARLIER THAT

$$\mathcal{L}(f(t)) = \frac{\int_0^T e^{-st} f(t) dt}{1 - e^{-sT}}$$

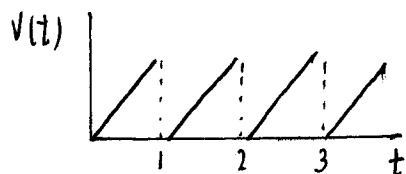
AS AN EXAMPLE OF THE USE OF THIS CONSIDER AN RL CIRCUIT WITH CURRENT $I(t)$ AS SHOWN



THEN $L \frac{dI}{dt} + RI = V(t) \quad I(0) = 0.$

$L, R > 0$ CONSTANTS

SUPPOSE THAT $V(t)$ IS A PERIODIC SAWTOOTH FUNCTION AS SHOWN



$V(t+1) = V(t)$ FOR $t > 1$

$V(t) = t$ FOR $0 < t < 1$

NOW LET $J(s) = \mathcal{L}(I(t))$ WE WANT TO DETERMINE $J(s)$ AND THEN RECOVER $I(t)$. WE WRITE

$$\mathcal{L}\left(L \frac{dI}{dt}\right) + R \mathcal{L}(I) = \mathcal{L}(V(t)) = V(s)$$

$$L[sJ - I(0)] + RJ = V(s) = \frac{\int_0^1 t e^{-st} dt}{1 - e^{-s}} = \frac{1}{1 - e^{-s}} \left[\left(-\frac{1}{s^2} - \frac{t}{s}\right) e^{-st} \Big|_0^1 \right]$$

HENCE $J(sL + R) = \frac{1}{1 - e^{-s}} \left(\left(-\frac{1}{s^2} - \frac{1}{s}\right) e^{-s} + \frac{1}{s^2} \right) = \frac{1}{(1 - e^{-s})} \left(+\frac{1}{s^2} (1 - e^{-s}) - \frac{1}{s} e^{-s} \right)$

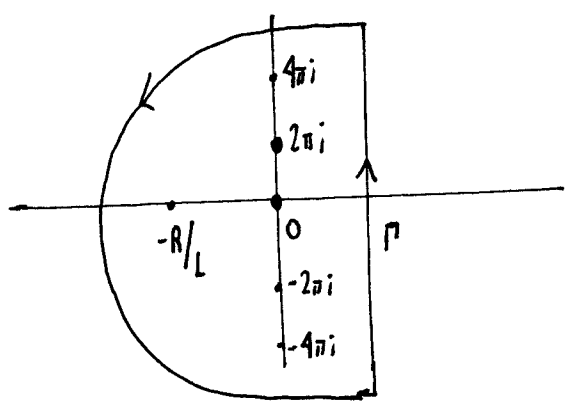
so $J(sL + R) = \frac{1}{s^2} - \frac{e^{-s}}{s(1 - e^{-s})} = \frac{1}{s^2} - \frac{1}{s(e^s - 1)} = \frac{1}{s^2} + \frac{1}{s(1 - e^s)}$

so $\mathcal{L}(I(t)) = J(s) = \left(\frac{1}{s^2} + \frac{1}{s(1 - e^{+s})} \right) \frac{1}{sL + R}$

NOTICE THAT WE HAVE A SIMPLE POLE AT $s = -R/L < 0$,

A DOUBLE POLE AT $s = 0$ AND SIMPLE POLES WHEN $1 - e^s = 0$ WHICH

GIVES $s = 2k\pi i$ $k = \dots, \pm 1, \pm 2, \dots$



$$\text{so } I(t) = \frac{1}{2\pi i} \int_{\Gamma} J(s) e^{st} ds$$

$$\begin{aligned} (*) \quad &= \text{REJ} [J e^{st}; 0] + \text{REJ} [J e^{st}; -R/L] \\ &+ \sum_{\substack{k=-\infty \\ k \neq 0}}^{\infty} \text{REJ} [e^{st} J; 2k\pi i] \end{aligned}$$

(i) NOW FOR $s = -R/L$ WE OBTAIN BY P/Q RESIDUE EVALUATION METHOD

THAT

$$\text{REJ} [J e^{st}; -R/L] = \frac{e^{-Rt/L}}{L} \left(\frac{1}{s^2} + \frac{1}{s(1-e^s)} \right) \Big|_{s=-R/L} = \frac{e^{-Rt/L}}{L} \left(\frac{L^2}{R^2} - \frac{L}{R(1-e^{-R/L})} \right)$$

$$\text{so } \text{RES} [J e^{st}; -R/L] = e^{-Rt/L} \left(\frac{L}{R^2} - \frac{1}{R(1-e^{-R/L})} \right)$$

(ii) FOR $s = 0$ WE HAVE A "DOUBLE" POLE. WE CALCULATE THE a_{-1} TERM FROM LAURENT SERIES

$$\text{REJ} [J(s) e^{st}; 0] \rightarrow \frac{1}{R} \left(\frac{1}{s^2} + \frac{1}{s(1-e^{-s})} \right) \frac{e^{st}}{1+sL/R}$$

$$\begin{aligned} \text{NEAR } s=0 \quad &\frac{1}{R} \left(\frac{1}{s^2} + \frac{1}{s(-\frac{s^2}{2}-s)} \right) \frac{e^{st}}{1+sL/R} \rightarrow \frac{1}{R s^2} \left(1 - \frac{1}{1+s/2} \right) (1+st+\dots)(1-sL/R) \\ &\rightarrow \frac{1}{R s^2} \left(1 - (1-s/2) \right) (1+st+\dots)(1-sL/R) \\ &\rightarrow \frac{1}{R s^2} \left(\frac{s}{2} \right) (1+\dots)(1-\dots) = \frac{1}{2RS} + \dots \end{aligned}$$

THEREFORE $s = 0$ IS ACTUALLY A POLE OF ORDER ONE AND

$$\text{REJ} [J(s) e^{st}; 0] = \frac{1}{2R}$$

(iii) NOW CALCULATE THE RESIDUE AT $s = 2\pi i k$ $k = \pm 1, \pm 2, \dots$

(20)

WE USE P/Q METHOD

$$J(s)e^{st} = \left(\frac{(1 - e^s) + s}{s^2(1 - e^s)} \right) \frac{e^{st}}{sL + R} = \frac{P}{Q} \quad Q = 1 - e^s$$

$$\text{NOW RES} [J e^{st}; 2\pi i k] = \frac{s e^{st}}{s^2(sL + R)(-e^s)} \Big|_{2\pi i k} = \frac{-e^{2\pi i k t}}{(2\pi i k)^2 L + 2\pi i k R}$$

THIS GIVES IN TOTAL FROM (X) THAT

$$I(t) = e^{-Rt/L} \left(\frac{L}{R^2} - \frac{1}{R(1 - e^{-R/L})} \right) + \frac{1}{2R} + \sum_{\substack{k=-\infty \\ k \neq 0}}^{\infty} \left(\frac{-e^{2\pi i k t}}{(2\pi i k)^2 L + 2\pi i k R} \right)$$

THIS CAN BE WRITTEN AS

$$I(t) = e^{-Rt/L} \left(\frac{L}{R^2} - \frac{1}{R(1 - e^{-R/L})} \right) + \frac{1}{2R} - \sum_{k=1}^{\infty} \left(\frac{e^{2\pi i k t}}{-4\pi^2 k^2 L + 2\pi i k R} + \frac{e^{-2\pi i k t}}{-4\pi^2 k^2 L - 2\pi i k R} \right)$$

$$\text{LET } z_k = \frac{e^{2\pi i k t}}{-4\pi^2 k^2 L + 2\pi i k R} \quad \text{THEN } \sum_{k=1}^{\infty} (z_k + \bar{z}_k) = 2 \sum_{k=1}^{\infty} \text{RE}(z_k)$$

$$\text{SO } I(t) = e^{-Rt/L} \left(\frac{L}{R^2} - \frac{1}{R(1 - e^{-R/L})} \right) + \frac{1}{2R} + 2 \sum_{k=1}^{\infty} \text{RE} \left[\frac{e^{2\pi i k t}}{-4\pi^2 k^2 L + 2\pi i k R} \right]$$

$$\text{NOW } \text{RE} \left[\frac{e^{2\pi i k t}}{-4\pi^2 k^2 L + 2\pi i k R} \right] = \text{RE} \left[\frac{(\cos(2\pi k t) + i \sin(2\pi k t))(-4\pi^2 k^2 L - 2\pi i k R)}{16\pi^4 k^4 L^2 + 4\pi^2 k^2 R^2} \right]$$

THIS YIELDS

$$I(t) = e^{-Rt/L} \left(\frac{L}{R^2} - \frac{1}{R(1 - e^{-R/L})} \right) + \frac{1}{2R} - 2 \sum_{k=1}^{\infty} \left(\frac{-4\pi^2 k^2 L \cos(2\pi k t) + 2\pi k R \sin(2\pi k t)}{16\pi^4 k^4 L^2 + 4\pi^2 k^2 R^2} \right)$$

NOTICE THAT THE SAWTOOTH FUNCTION GIVES RISE TO INFINITE # OF MODES OF OSCILLATION.