

On a Semilinear Elliptic Equation in R^2 When the Exponent Approaches Infinity

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We consider a semilinear elliptic equation in R^2 with the nonlinear exponent approaching infinity. In contrast to the blow-up behavior of the corresponding problem in R^n with $n \geq 3$, the $L^\infty(R^2)$ norms of the solutions to the equation in R^2 remain bounded from below and above. After a careful study on the decay rates of several quantities, we prove that the normalized solutions will approach the fundamental solution of $-\Delta u + 1$ in R^2 . So as the exponent tends to infinity, the solutions to the problem look more and more like a peak. © 1995 Academic Press, Inc.

1. INTRODUCTION

This paper is devoted to the behavior of the ground state solution to a semilinear elliptic equation with large exponent of the nonlinear term. Considering

$$\begin{cases} \Delta u - u + u^p = 0 & \text{in } R^2 \\ u > 0, \quad \lim_{|x| \rightarrow \infty} u(x) = 0, \end{cases} \quad (1.1)$$

we would like to understand the behavior of the solutions to (1.1) when the exponent p approaches ∞ .

The corresponding problem in higher dimensions

$$\begin{cases} \Delta u - u + u^p = 0 & \text{in } R^n, \text{ with } n \geq 3 \\ u > 0, \quad \lim_{|x| \rightarrow \infty} u = 0, \end{cases} \quad (1.2)$$

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was studied by X. Pan and X. Wang in [9]. Due to the existence of the critical exponent, namely $(n+2)/(n-2)$ there, they were led to study the behavior of the solution to the corresponding equation in higher dimensions when p approaches $(n+2)/(n-2)$. It turns out that solutions to their problem will blow up as $p \rightarrow (n+2)/(n-2)$.

Because of the critical exponent for $n \geq 3$, in higher dimensions we have a profile equation

$$\begin{cases} \Delta u - u + u^{(n+2)/(n-2)} = 0 & \text{in } R^n \\ u > 0, \quad \lim_{|x| \rightarrow \infty} u(x) = 0. \end{cases} \quad (1.3)$$

In fact the nonexistence of solutions to (1.3) implies the blow-up of solutions to (1.2) when $p \rightarrow (n+2)/(n-2)$.

In R^2 , however, due to the lack of Eq. (1.3), we will approach the problem in a different way. We will start with a sharp estimate on the growth rate of c_p where c_p is the best constant of the embedding

$$W^{1,2}(R^2) \hookrightarrow L^{p+1}(R^2).$$

From this estimate and estimates of some other quantities, we will prove

THEOREM 1.1. *Let u_p be a solution to (1.1). Then for p large enough, there exists C independent of p such that $1 < \|u_p\|_{L^2(R^2)} < C$.*

THEOREM 1.2. *Let u_p be a radially symmetric solution to (1.1). Then, as $p \rightarrow \infty$, we have*

$$\begin{aligned} (1) \quad & \frac{u_p}{\int_{R^2} u_p} \rightarrow G \quad \text{in the sense of distribution,} \\ (2) \quad & \frac{u_p}{\int_{R^2} u_p} \rightarrow G \quad \text{in } C_{loc}^{2,\alpha}(R^2 \setminus \{0\}) \text{ for all } \alpha \in (0, 1), \end{aligned}$$

where G is the fundamental solution to $-\Delta + 1$, i.e.,

$$-\Delta G + G = \delta.$$

Remark 1.3. G has an integral representation

$$G(x) = \frac{1}{\Gamma(1/2)} \sqrt{\frac{\pi}{2}} \frac{e^{-x}}{\sqrt{x}} \int_0^\infty e^{-t} t^{-1/2} \left(1 + \frac{1}{2} \frac{t}{x}\right)^{-1/2} dt.$$

Moreover, G decays exponentially at infinity. We refer to [7] for more information.

We see, in contrast to the higher dimensions, the solutions look more and more like a peak as p gets large.

Our paper is organized as follows. In Section 2, we give necessary background of the solutions to (1.1) together with some integral identities. Then we prove a growth rate estimate of c_p in Section 3. Finally we prove Theorem 1.1 in Section 4 and Theorem 1.2 in Section 5.

2. PRELIMINARIES

There are various ways to obtain solutions to (1.1). We shall adapt the variational approach developed in [2].

Let

$$\mathcal{A}_p = \left\{ u \in W^{1,2}(R^2) : \int_{R^2} u^{p+1} = 1 \right\}.$$

Consider the functional $J_p: \mathcal{A}_p \rightarrow R$ defined by

$$J_p(u) = \frac{1}{2} \int_{R^2} [|\nabla u|^2 + u^2].$$

By the theory of Schwartz symmetrization (see, for example, [6]), we can replace any minimizing sequence of the functional by a radially symmetric minimizing sequence. Then it is easy to show, with the aid of a compactness result of radial embedding in [10], that J_p has a radial minimizer in \mathcal{A}_p . If we denote this minimizer by u'_p , a scalar multiple of u'_p will solve Eq. (1.1). We denote this radial solution by u_p . A further L^∞ estimate shows $\lim_{|x| \rightarrow \infty} u(x) = 0$. Hence u_p solves (1.1).

According to a result in [3] all solutions of (1.1) up to a parallel translation are radially symmetric. Applying the uniqueness result in [8] or [5], we know that the radial solution is unique. Therefore (1.1) has a unique solution up to a translation. From now on we will denote this unique radial solution by u_p .

It is easy to see that u_p is related to the embedding

$$W^{1,2}(R^2) \hookrightarrow L^{p+1}(R^2).$$

We refer to [1] for the proof of this embedding theorem.

If we denote c_p as the best constant of the above embedding, i.e., c_p is the least number among all possible constant C 's which make the inequality

$$\|u\|_{L^{p+1}(R^2)} \leq C \|u\|_{W^{1,2}(R^2)}$$

for all $u \in W^{1,2}(R^2)$, then c_p is achieved by u_p . Hence we have

LEMMA 2.1. *Let u_p be the radial solution to (1.1) as above. Then we have*

$$\frac{\|u_p\|_{W^{1,2}(R^2)}}{\|u_p\|_{L^{p+1}(R^2)}} = \inf \left\{ \frac{\|u\|_{W^{1,2}(R^2)}}{\|u\|_{L^{p+1}(R^2)}} : u \in W^{1,2}(R^2), u \neq 0 \right\} = \frac{1}{c_p},$$

where c_p is the best constant defined above.

Now let us state some integral identities.

LEMMA 2.2. *Let u_p be a solution to (1.1). We have*

$$\begin{aligned} (1) \quad & \int_{R^2} [|\nabla u_p|^2 + |u_p|^2] = \int_{R^2} u_p^{p+1} \\ (2) \quad & \int_{R^2} u_p = \int_{R^2} u_p^p \\ (3) \quad & \frac{1}{2} \int_{R^2} u_p^2 = (1/(p+1)) \int_{R^2} u_p^{p+1}. \end{aligned}$$

Proof. (1) Multiplying (1.1) by u_p and integrating over $B(R)$, the ball of radius R centered at the origin, we get

$$\int_{B(R)} [|\nabla u_p|^2 + |u_p|^2] = \int_{\partial B_R} u_p \frac{\partial u_p}{\partial n} + \int_{B_R} u_p^{p+1}.$$

Using an exponential decay property of u_p which says that there exist C, μ independent of x such that

$$u_p(x) \leq C e^{-\mu|x|}, \quad |Du_p(x)| \leq C e^{-\mu|x|}$$

we can let R approach ∞ yielding 1. The proof of the exponential decay estimate can be found, for example, in [2]. We shall prove a refined result, Lemma 5.1, by the same method in Section 5.

(2) Similar to (1), but we integrate the equation directly.

(3) This is the R^2 version of the well known Pohozaev identity. We refer to [2] for a proof. ■

We also present some simple radial lemmas which will be used in later sections.

LEMMA 2.3. *Let $n \geq 2$. Every radial function $u \in W^{1,2}(R^n)$ is almost equal to a function $U(x)$ which is continuous for $x \neq 0$; furthermore*

$$|U(x)| \leq C_n |x|^{(1-n)/2} \|u\|_{W^{1,2}(R^n)}$$

for $|x| \geq \alpha_n$ where C_n, α_n depend on n only.

The proof of this lemma can be found in [2].

LEMMA 2.4. *Let $u \in L^2(R^2)$ be a radial, nonnegative, non-increasing in $|x| > 0$, continuous function except at 0. Then we have*

$$u(|x|) \leq \frac{2}{\sqrt{3\pi}} |x|^{-1} \|u\|_{L^2(R^2)}.$$

Proof.

$$\begin{aligned} \|u\|_{L^2(R^2)}^2 &= 2\pi \int_0^\infty u^2(r)r \, dr \geq 2\pi \int_{|x|/2}^{|x|} u^2(s)s^{-1} \, ds \\ &\geq 2\pi \frac{1}{2} u^2(|x|) \left(|x|^2 - \frac{|x|^2}{4} \right) \geq \frac{3\pi}{4} |x|^2 u^2(|x|). \end{aligned}$$

Therefore we have the lemma. ■

3. ON THE GROWTH RATE OF c_p

In this section, we establish a sharp estimate for the growth rate of c_p . Notice that according to the Schwartz symmetrization theory, the c_p is actually obtained in the class of radial functions. We will use C to denote various constants independent of p .

LEMMA 3.1. *If $u \in W_0^{1,2}(\Omega)$ where Ω is a bounded smooth domain in R^2 , then for every $t \geq 2$*

$$\|u\|_{L^t(\Omega)} \leq Ct^{1/2} |\Omega|^{1/t} \|\nabla u\|_{L^2(\Omega)}.$$

Proof. Let $u \in W_0^{1,2}(\Omega)$. We know

$$\frac{1}{\Gamma(s+1)} x^s \leq e^x$$

for all $x \geq 0, s \geq 0$ where Γ is the Γ function. From Trudinger's Inequality (see [4, p. 160]), we have

$$\int_{\Omega} \exp \left[C_1 \left(\frac{u}{\|\nabla u\|_{L^2}} \right)^2 \right] dx \leq C_2 |\Omega|,$$

where C_1, C_2 are constants which depend on the dimension of Ω only and $|\cdot|$ denotes the Lebesgue measure. Therefore

$$\begin{aligned} & \frac{1}{\Gamma(t/2 + 1)} \int_{\Omega} u^t dx \\ &= \frac{1}{\Gamma(t/2 + 1)} \int_{\Omega} \left[C_1 \left(\frac{u}{\|\nabla u\|_{L^2}} \right)^2 \right]^{t/2} dx C_1^{-t/2} \|\nabla u\|_{L^2}^t \\ &\leq \int_{\Omega} \exp \left[C_1 \left(\frac{u}{\|\nabla u\|_{L^2}} \right)^2 \right] dx C_1^{-t/2} \|\nabla u\|_{L^2}^t \leq C_2 |\Omega| C_1^{-t/2} \|\nabla u\|_{L^2}^t. \end{aligned}$$

Hence

$$\left(\int_{\Omega} u^t dx \right)^{1/t} \leq \left(\Gamma \left(\frac{t}{2} + 1 \right) \right)^{1/t} C_2^{1/t} C_1^{-1/2} |\Omega|^{1/t} \|\nabla u\|_{L^2(\Omega)}.$$

Notice that, according to Stirling's formula,

$$\left(\Gamma \left(\frac{t}{2} + 1 \right) \right)^{1/t} \sim \left(\left(\frac{t/2}{e} \right)^{t/2} \sqrt{t e} e^{\theta_t} \right)^{1/t} \sim C t^{1/2},$$

where $0 < \theta_t < 1/12$. This completes our lemma. \blacksquare

LEMMA 3.2. For each $u \in W^{1,2}(R^2)$, we have

$$\|u\|_{L^t(R^2)} \leq C t^{1/2} \|u\|_{W^{1,2}(R^2)}$$

for every $t \geq 4$.

Proof. As we mentioned earlier, by the Schwartz symmetrization theory, we need only to prove the inequality for radial functions. So in the proof all functions are assumed to be radial. Let $u := u_1 + u_2 := u \chi_{\{r \leq \alpha_2\}} + u \chi_{\{r > \alpha_2\}}$ where α_2 is defined in Lemma 2.3.

Then

$$\|u\|_{L^t(R^2)} \leq \|u_1\|_{L^t} + \|u_2\|_{L^t}.$$

And by Lemma 2.3

$$\begin{aligned} \|u_2\|_{L^t} &\leq \left(\int_{R^2 \setminus B_{\alpha_2}} u^t \right)^{1/t} \leq \left(2\pi \int_{r > \alpha_2} r^{-t/2+1} (C \|u\|_{W^{1,2}(R^2)})^t \right)^{1/t} \\ &\leq C \|u\|_{W^{1,2}(R^2)} \left(\frac{2}{|t-1|} \right)^{1/t} (\alpha_2^{-t/2+2})^{1/t} \leq C \|u\|_{W^{1,2}(R^2)}. \end{aligned}$$

Hence combining this with Lemma 3.1, we obtain

$$\|u\|_{L^1} \leq Ct^{1/2}\|u\|_{W^{1,2}(R^2)} + C\|u_2\|_{W^{1,2}(R^2)} \leq Ct^{1/2}\|u\|_{W^{1,2}(R^2)}. \quad \blacksquare$$

LEMMA 3.3. *For p large enough, there exist C_1, C_2 independent of p such that*

$$C_1p^{1/2} \leq c_p \leq C_2p^{1/2}.$$

Proof. The second inequality follows from Lemma 3.2. To prove the first inequality, we consider the so-called Moser's function

$$\begin{aligned} \Phi(r) &= \begin{cases} (p+1)^{-1/2} \log \frac{1}{r}, & e^{-(p+1)} \leq r \leq 1 \\ (p+1)^{-1/2} \log e^{p+1}, & 0 \leq r \leq e^{-(p+1)} \end{cases} \\ &= \begin{cases} (p+1)^{-1/2} \log \frac{1}{r}, & e^{-(p+1)} \leq r \leq 1 \\ (p+1)^{1/2}, & 0 \leq r \leq e^{-(p+1)}. \end{cases} \end{aligned}$$

So $\phi \in W_0^{1,2}(B_1) \subset W^{1,2}(R^2)$ where the trivial extension is used to interpret the second inclusion.

On the one hand

$$\|\nabla\phi\|_{L^2}^2 = 2\pi \int_0^1 \phi_r^2 r \, dr = 2\pi \int_{e^{-(p+1)}}^1 (p+1)^{-1} \frac{1}{r} \, dr = 2\pi.$$

Therefore by Poincaré's Inequality on B_1 ,

$$\|\phi\|_{W^{1,2}(R^2)} = \|\phi\|_{W^{1,2}(B_1)} \leq C\|\nabla\phi\|_{L^2(B_1)} = C, \quad (3.1)$$

where C is a constant depending on the first eigenvalue of Δ on the unit disk with Dirichlet boundary condition only.

On the other hand

$$\begin{aligned} \int_{B_1} \phi^{p+1} &= 2\pi \int_0^1 \phi^{p+1} r \, dr \geq 2\pi \int_{e^{-(p+1)}}^1 (p+1)^{-(p+1)/2} \left(\log \frac{1}{r}\right)^{p+1} r \, dr \\ &= 2\pi(p+1)^{-(p+1)/2} \int_0^{p+1} t^{p+1} e^{-t} \, dt \\ &\geq 2\pi(p+1)^{-(p+1)/2} \int_{(p+1)/2}^{p+1} t^{p+1} e^{-t} \, dt \\ &\geq 2\pi(p+1)^{-(p+1)/2} \left(\frac{p+1}{2}\right)^{p+1} e^{-(p+1)} \frac{p+1}{2} \\ &\geq 2\pi 2^{-(p+2)} (p+1)^{(p+3)/2} e^{-(p+1)}. \end{aligned}$$

Hence

$$\|\phi\|_{L^{p+1}} \geq (2\pi)^{1(p+1)} 2^{-(p+2)/(p+1)} (p+1)^{(1/2)((p+3)/(p+1))} e^{-1} \geq C(p+1)^{1/2}. \quad (3.2)$$

Combining (3.1) and (3.2), we have the lemma. \blacksquare

4. ON THE $L^\infty(\mathbb{R}^2)$ NORM OF u_p

We are now in the position to study the $L^\infty(\mathbb{R}^2)$ norms of u_p . Because our solutions u_p are radially symmetric, they satisfy the following ordinary differential equation

$$\begin{cases} u'' + \frac{1}{r} u' - u + u^p = 0 \\ u'(0) = 0, \quad \lim_{r \rightarrow \infty} u(r) = 0. \end{cases} \quad (4.1)$$

From [8] or [5], we know $u'_p(r) < 0$ for all $r > 0$ where u_p is a solution to (1.1) with the exponent in the equation being p . We define

$$\gamma_p := u_p(0) = \|u_p\|_{L^\infty}.$$

LEMMA 4.1. $\|u_p\|_{W^{1,2}(\mathbb{R}^2)} \leq C/p^{1/2}$ when p is large.

Proof. From the integral identity (1) in Lemma 2.2, we deduce

$$\|u_p\|_{W^{1,2}(\mathbb{R}^2)}^2 = \|u_p\|_{L^{p+1}(\mathbb{R}^2)}^{p+1}.$$

Combining this with Lemma 2.1 and Lemma 3.2, we get

$$\|u_p\|_{W^{1,2}(\mathbb{R}^2)} = C^{(p+1)/(p-1)} p^{-(1/2)(p+1)/(p-1)} \leq Cp^{-1/2}$$

when p is large. \blacksquare

We now prove Theorem 1.1.

Proof of Theorem 1.1. The uniform lower bound 1 follows from the maximum principle. To find a uniform upper bound let

$$E_p(r) = -ru'_p(r) - \frac{1}{2} r^2(u_p^p - u).$$

Then $E_p(0) = 0$, and

$$\begin{aligned} E_p'(r) &= -u_p' - ru_p'' - r(u_p^p - u_p) - \frac{1}{2} r^2 (pu_p^{p-1} - 1)u_p' \\ &= -\frac{1}{2} r^2 (pu_p^{p-1} - 1)u_p'. \end{aligned}$$

Let $r_0 > 0$ be such that $u_p(r_0) = (1/p)^{1/(p-1)}$. Clearly such r_0 exists, and it depends on p . Notice $E_p'(r) \geq 0$ for $r \in (0, r_0)$. Therefore $E_p \geq 0$ on $(0, r_0)$. From the definition of E , we have

$$-u_p'(r) - \frac{1}{2} r(u_p^p - u_p) \geq 0$$

on $(0, r_0)$. Combining this with Eq. (4.1) of u_p , we get

$$u_p'' \geq \frac{1}{2} (u_p - u_p^p), \quad u_p'' \geq \frac{1}{2} (\gamma_p - \gamma_p^p) \tag{4.2}$$

for all r in $(0, r_0)$.

Integrating both sides of (4.2) twice with respect to r on $(0, r)$, we obtain

$$\gamma_p - \frac{1}{2} u_p(r) \leq \frac{1}{4} r^2 (\gamma_p^p - \gamma_p) \tag{4.3}$$

for $r \in (0, r_0)$.

Now let T_0 be such that $u_p(T_0) = \frac{1}{2}\gamma_p$. T_0 depends on p . From (4.3), we have

$$\gamma_p - \frac{1}{2} \gamma_p \leq \frac{1}{4} T_0^2 (\gamma_p^p - \gamma_p), \quad \gamma_p \leq \frac{1}{2} T_0^2 \gamma_p^p. \tag{4.4}$$

Applying Lemma 3.2 with $t = 2p$ and Lemma 4.1, we get

$$\|u_p\|_{L^{2p}(R^2)} \leq C(2p)^{1/2} \|u_p\|_{W^{1,2}(R^2)} \leq C(2p)^{1/2} C' p^{-1/2} := M,$$

where M is independent of p for large p . Hence

$$\int_{R^2} u_p^{2p} dx \leq M^{2p}.$$

But

$$\int_{R^2} u^{2p} dx = 2\pi \int_0^\infty u_p^{2p}(r)r dr \geq \int_0^{T_0} u_p^{2p} r dr \geq \left(\frac{\gamma_p}{2}\right)^{2p} \frac{1}{2} T_0^2.$$

Therefore

$$\frac{1}{2} T_0^2 \left(\frac{\gamma_p}{2}\right)^{2p} \leq M^{2p}, \quad T_0^2 \leq 2 \left(\frac{2M}{\gamma_p}\right)^{2p}. \quad (4.5)$$

Combining (4.4) with (4.5), we obtain

$$1 \leq \left(\frac{2M}{\gamma_p}\right)^{2p} (\gamma_p^{p-1}) \leq \left(\frac{2M}{\gamma_p}\right)^{2p} \gamma_p^p, \quad \gamma_p \leq (2M)^2. \quad \blacksquare$$

We now derive the decay rates for some quantities.

COROLLARY 4.2. *There exist C_1, C_2 independent of p such that for large p we have*

$$(1) \quad C_1 \frac{1}{p} \leq \int_{R^2} |\nabla u_p|^2 dx \leq C_2 \frac{1}{p}$$

$$(2) \quad C_1 \frac{1}{p} \leq \int_{R^2} u_p^{p+1} dx \leq C_2 \frac{1}{p}$$

$$(3) \quad C_1 \frac{1}{p^2} \leq \int_{R^2} u_p^2 dx \leq C_2 \frac{1}{p^2}$$

$$(4) \quad C_1 \frac{1}{p} \leq \int_{R^2} u_p dx = \int_{R^2} u_p^p dx \leq C_2 \frac{1}{p}.$$

Proof. Combining Lemma 2.1, Lemma 2.2, and Lemma 3.3 we have (1)–(3). To prove (4), we see from Theorem 1.1 and (2)

$$\int_{R^2} u_p^p \geq \frac{1}{\gamma_p} \int_{R^2} u_p^{p+1} \geq \frac{1}{C} \int_{R^2} u_p^{p+1} \geq \frac{C_1}{p}.$$

On the other hand, an interpolation argument shows

$$\|u_p\|_{L^p} \leq \|u_p\|_{L^2}^{(p-2)/(p+1-2)} \|u_p\|_{p+1}^{(p+1-p)/(p+1-2)}.$$

Therefore (2) and (3) imply

$$\int_{R^2} u_p^p \leq \frac{C_2}{p}$$

for C_2 independent of p when p is large. Therefore with the aid of Lemma 2.2(2) we have (4). ■

5. ON THE NORMALIZED u_p

Before we turn to the proof of Theorem 1.2, we first give a uniform exponential decay estimate in p of u_p .

LEMMA 5.1. *There exists a constant C independent of p , such that for large p and $r \geq 1$*

$$u_p(r) \leq \frac{C}{p} e^{-r/2} r^{-1/2}.$$

Proof. Let $v_p = r^{1/2} u_p$. Then v_p satisfies

$$v_p'' = \left[q(r) - \frac{1}{4r^2} \right] v_p,$$

where $q(r) = 1 - u_p^{p-1}(r)$. Note for large p , by Lemma 2.4 and Corollary 4.2

$$q(1) - \frac{1}{41^2} \geq 1 - \left[\frac{2}{\sqrt{3\pi}} \|u_p\|_{L^2} \right]^{p-1} - \frac{1}{4} \geq \frac{3}{4} - \left(\frac{C}{p} \right)^{p-1} \rightarrow \frac{3}{4}$$

as $p \rightarrow \infty$. Hence for large p we have

$$q(r) - \frac{1}{4r^2} \geq \frac{1}{2}$$

for $r \geq 1$.

Let $w_p = v_p^2$; then w_p satisfies

$$\frac{1}{2} w_p'' = (v_p')^2 + \left[q(r) - \frac{1}{4r^2} \right] w_p.$$

Thus for $r \geq 1$, one has $w_p'' \geq w_p$ and $w_p \geq 0$.

Let $z_p = e^{-r}(w_p' + w_p)$. We have $z_p' = e^{-r}(w_p'' - w_p) \geq 0$; hence z_p is a nondecreasing function on $(1, \infty)$. If there exists $r_1 > 1$ such that $z_p(r_1) > 0$, then $z_p(r) \geq z_p(r_1) > 0$ for $r > r_1$. This implies that

$$w_p' + w_p \geq (z_p(r_1))e^r$$

whence $w'_p + w_p$ is not integrable on (r_1, ∞) . But v_p^2 and $v_p v'_p$ are integrable near ∞ for $u_p \in W^{1,2}(\mathbb{R}^2)$, so that w'_p and w_p are also integrable, a contradiction.

Hence $z_p(r) \leq 0$ for $r \geq 1$. This implies that

$$(e^r w_p)' = e^{2r} z_p \leq 0$$

for $r \geq 1$. Hence

$$w_p \leq e^1 w_p(1) e^{-r}, \quad u_p \leq e^{1/2} u_p(1) e^{-r/2} r^{-1/2}.$$

But $u_p(1)$ can be estimated by Lemma 2.4 and Corollary 4.2 as

$$u_p(1) \leq \frac{2}{\sqrt{3\pi}} \|u_p\|_{L^2(\mathbb{R}^2)} \leq \frac{C}{p}.$$

Therefore

$$u_p(r) \leq \frac{C}{p} e^{-r/2} r^{-1/2}$$

for $r \geq 1$. ■

Proof of Theorem 1.2. Let

$$A_\varepsilon = \left\{ x \in \mathbb{R}^2 : \varepsilon \leq |x| \leq \frac{1}{\varepsilon} \right\}.$$

Then by Lemma 2.4 and Corollary 4.2 we have that for large p

$$\left\| \frac{u_p}{\int u_p} \right\|_{L^2(A_\varepsilon)} \leq C_\varepsilon p \|u_p\|_{L^2} \leq C_\varepsilon.$$

Hence $\|u_p / \int u_p\|_{L^2(A_\varepsilon)}$ is bounded uniformly for large p .

On the other hand, by Lemma 2.4 and Corollary 4.2

$$\left\| \frac{u_p^p}{\int u_p^p} \right\|_{L^2(A_\varepsilon)} \leq Cp \left(\frac{C_2}{p} \right)^p \varepsilon^{-p}$$

which is bounded uniformly for large p . So the elliptic L^p -estimate shows that $u_p / \int u_p$ is bounded in $C^\alpha(A_\varepsilon)$ for all $\alpha \in (0, 1)$. Using the Schauder estimate, we have that $u_p / \int u_p$ is bounded uniformly in $C^{2,\alpha}(A_\varepsilon)$ for all $\alpha \in$

(0, 1). Therefore a subsequence of $\{u_p/\int u_p\}$, still denoted by $\{u_p/\int u_p\}$, will approach a function, say G' , in $C_{loc}^{2,\alpha}(R^2 \setminus \{0\})$. So we almost have (1) except that we need to show $G' = G$ where G is the fundamental solution given in Remark 1.3.

We now prove

$$\frac{u_p}{\int u_p} \rightarrow G'$$

in the sense of distribution and

$$-\Delta G' + G' = \delta.$$

Let $\phi(x) \in C_0^\infty(R^2)$, $\varepsilon > 0$. A standard kernel argument shows that (note $\int u_p^p = \int u_p$ by Lemma 2.2), with the aid of Lemma 2.4

$$\begin{aligned} \left| \int_{R^2} \phi(x) \frac{u_p^p}{\int u_p} dx - \phi(0) \right| &\leq \int_{R^2} |\phi(x) - \phi(0)| \frac{u_p^p}{\int u_p} \\ &\leq \int_{B_\delta} |\phi(x) - \phi(0)| \frac{u_p^p}{\int u_p} + 2 \max_{x \in R^2} (\phi(x)) \int_{R^2 \setminus B_\delta} \frac{u_p^p}{\int u_p} \\ &\leq \max_{B_\delta} (\phi(x) - \phi(0)) + 2 \max_{R^2} [\phi(x)] C_p \int_\delta^\infty \frac{1}{r^p} r \left(\frac{C}{p}\right)^p \\ &\leq \frac{\varepsilon}{2} + \frac{\varepsilon}{2} = \varepsilon \end{aligned}$$

if we choose δ small enough first, then choose p large enough.

Hence,

$$\lim_{p \rightarrow \infty} \int \phi(x) \frac{u_p^p}{\int u_p} = \phi(0).$$

But

$$\int \phi(x) \frac{u_p^p}{\int u_p} = \int [(-\Delta + 1)\phi(x)] \frac{u_p}{\int u_p}. \quad (5.1)$$

To pass the limit in Eq. (5.1), we need a dominating function.

Let us first consider $u_p/\int u_p$ in B_1 , the unit disc. Lemma 2.4 and Corol-

lary 4.2 imply that

$$\frac{u_p}{\int u_p} \leq \frac{2C}{|x|}. \quad (5.2)$$

So we take the latter to be the dominating function in B_1 .

Then we use Lemma 5.1 to get the following decay estimate on $R^2 \setminus B_1$

$$\frac{u_p(x)}{\int u_p} \leq C e^{-|x|/2} |x|^{-|x|/2}, \quad (5.3)$$

where the latter gives a dominating function on $R^2 \setminus B_1$.

Combining (5.2) with (5.3), we can pass the limit in (5.1); hence

$$\begin{aligned} \int [(-\Delta + 1)\phi(x)]G'(x) &= \lim_{p \rightarrow \infty} \int [(-\Delta + 1)\phi(x)] \frac{u_p(x)}{\int u_p} \\ &= \lim_{p \rightarrow \infty} \int \phi(x) \frac{u_p(x)}{\int u_p} = \phi(0). \end{aligned}$$

Therefore G' satisfies

$$-\Delta G' + G' = \delta;$$

hence

$$-\Delta(G - G') + (G - G') = 0,$$

where G is the fundamental solution given in Remark 1.3. Because $G - G'$ decays, actually decays exponentially by Remark 1.3 and (5.3), at infinity, the maximum principle implies $G = G'$; hence we get (1) and (2) simultaneously. ■

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