Lecture 29: The heat equation with Robin BC

(Compiled 4 August 2017)

In this lecture we demonstrate the use of the Sturm-Liouville eigenfunctions in the solution of the heat equation. We first discuss the expansion of an arbitrary function f(x) in terms of the eigenfunctions $\{\phi_n(x)\}$ associated with the Robins boundary conditions. This is a generalization of the Fourier Series approach and entails establishing the appropriate normalizing factors for these eigenfunctions. We then uses the new generalized Fourier Series to determine a solution to the heat equation when subject to Robins boundary conditions.

Key Concepts: Eigenvalue Problems, Sturm-Liouville Boundary Value Problems; Robin Boundary conditions.

Reference Section: Boyce and Di Prima Section 11.1 and 11.2

29 Solving the heat equation with Robin BC

29.1 Expansion in Robin Eigenfunctions

In this subsection we consider a Robin problem in which $\ell = 1$, $\mathbf{h_1} \to \infty$, and $\mathbf{h_2} = \mathbf{1}$, which is a Case III problem as considered in lecture 30. In particular:

$$\begin{array}{c} \phi'' + \mu^2 \phi = 0\\ \phi(0) = 0, \phi'(1) = -\phi(1) \end{array} \right\} \Longrightarrow \left\{ \begin{array}{c} \phi_n = \sin(\mu_n x),\\ \tan(\mu_n) = -\mu_n\\ \mu_n \sim \left[\left(\frac{2n+1}{2}\right) \pi \right] \text{ as } n \to \infty \end{array} \right.$$

Assume that we can expand f(x) in terms of $\phi_n(x)$:

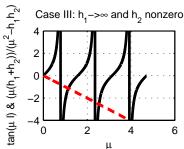
$$f(x) = \sum_{n=1}^{\infty} c_n \phi_n(x) \tag{29.1}$$

$$\int_{0}^{1} f(x)\sin(\mu_{n}x) \, dx = c_{n} \int_{0}^{1} \left[\phi_{n}(x)\right]^{2} \, dx \tag{29.2}$$

$$= c_n \frac{1}{2} \left[1 + \cos^2 \mu_n \right]$$
 (29.3)

Therefore

$$c_n = \frac{2}{\left[1 + \cos^2 \mu_n\right]} \int_0^1 f(x) \sin(\mu_n x) \, dx.$$
(29.4)



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If f(x) = x then

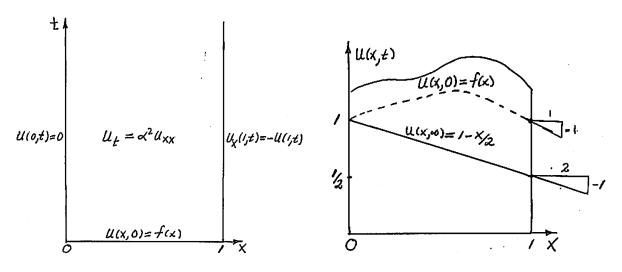
but

$$\int_{0}^{1} x \sin(\mu_n x) dx = -\frac{\cos(\mu_n x)}{\mu_n} - x \Big|_{0}^{1} + \frac{1}{\mu_n} \int_{0}^{1} \cos \mu_n x dx$$
$$= -\frac{\cos(\mu_n)}{\mu_n} + \frac{\sin \mu_n x}{\mu_n^2} \Big|_{0}^{1}$$
(29.5)
$$= \frac{\sin \mu_n - \mu_n \cos \mu_n}{\mu_n^2} = 2 \frac{\sin \mu_n}{\mu_n^2}.$$

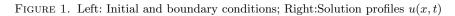
Therefore

$$c_n = \frac{4\sin\mu_n}{\mu_n^2 [1 + \cos^2\mu_n]}$$
(29.6)
$$f(x) = 4\sum_{n=1}^{\infty} \frac{\sin\mu_n \sin(\mu_n x)}{\mu_n^2 [1 + \cos^2\mu_n]}$$
(29.7)

29.2 Solving the Heat Equation with Robin BC



(b) Solution profiles u(x,t) at various times



$$u_t = \alpha^2 u_{xx} \quad 0 < x < 1 \tag{29.8}$$

$$u(0,t) = 1 \quad u_x(1,t) + u(1,t) = 0 \tag{29.9}$$

$$u(x,0) = f(x).$$
(29.10)

Look for a steady state solution v(x)

$$\begin{cases} v''(x) = 0 \\ v(0) = 1 \quad v'(1) + v(1) = 0 \end{cases}$$
 (29.11)

$$v = Ax + B \quad v(0) = B = 1 \quad v'(x) = A \quad v'(1) + v(1) = A + (A+1) = 0$$

$$A = -1/2 \tag{29.12}$$

Therefore

$$v(x) = 1 - x/2. \tag{29.13}$$

Now let u(x,t) = v(x) + w(x,t)

$$u_t = w_t = \alpha^2 (v'' + w_{xx}) \Rightarrow w_t = \alpha^2 w_{xx}$$
$$1 = u(0, t) = v(0) + w(0, t) = 1 + w(0, t) \Rightarrow w(0, t) = 0$$
$$(1, t) + w(1, t) = -(w'(1) + 2w(1)) = + w(1, t) + w(1, t) = -w(1, t) + w(1, t)$$

$$\begin{array}{rcl} 0 = u_x(1,t) + u(1,t) &=& \{v'(1) \not \to v(1)\} &+ w_x(1,t) + w(1,t) &\Rightarrow& w_x(1,t) + w(1,t) = 0 \\ f(x) = u(x,0) &=& v(x) + w(x,0) &\Rightarrow& w(x,0) = f(x) - v(x). \end{array}$$

Let

$$w(x,t) = X(x)T(t)$$
 (29.14)

$$\frac{T(t)}{\alpha^2 T(t)} = \frac{X''}{X} = -\mu^2$$
(29.15)

$$T(t) = c e^{-\alpha^2 \mu^2 t}$$
(29.16)

$$X'' + \mu^2 X = 0
 X(0) = 0 \quad X'(1) + X(1) = 0
 \begin{cases}
 The \ \mu_n \ \text{are solutions of the transcendental} \\
 equation: \ \tan \mu_n = -\mu_n.
 \end{cases}$$
(29.17)

$$X_n(x) = \sin(\mu_n x) \tag{29.18}$$

$$w(x,t) = \sum_{n=1}^{\infty} c_n e^{-\alpha^2 \mu_n^2 t} \sin(\mu_n x)$$
(29.19)

where

$$f(x) - v(x) = w(x,0) = \sum_{n=1}^{\infty} c_n \sin(\mu_n x)$$
(29.20)

$$\Rightarrow c_n = \frac{2}{[1 + \cos^2 \mu_n]} \int_0^1 [f(x) - v(x)] \sin(\mu_n x) \, dx \tag{29.21}$$

$$u(x,t) = 1 - \frac{x}{2} + \sum_{n=1}^{\infty} c_n e^{-\alpha^2 \mu_n^2 t} \sin(\mu_n x).$$
(29.22)