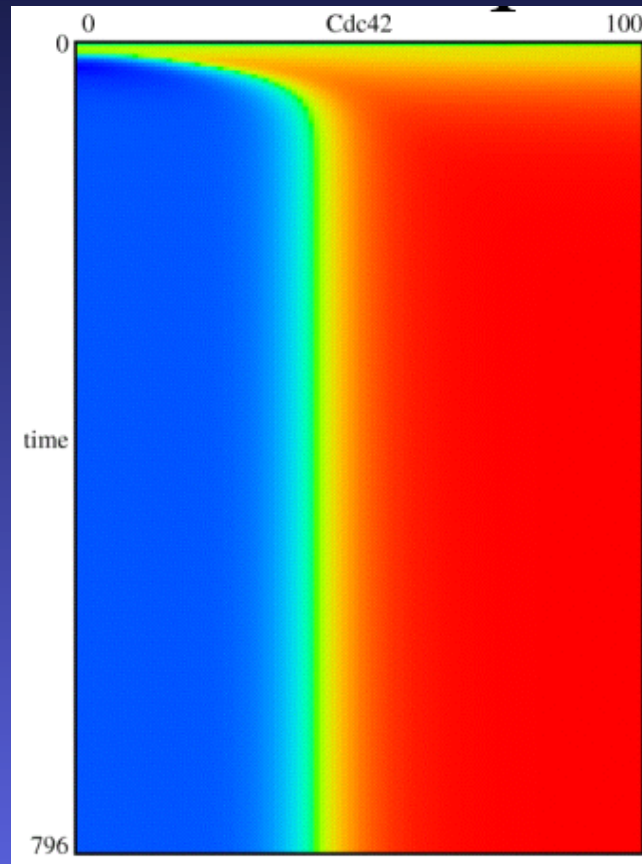
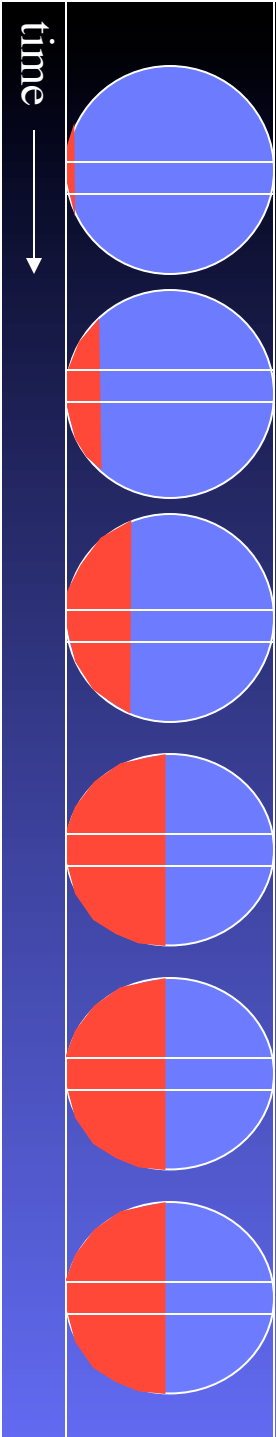


Mathematical Cell Biology Graduate Summer Course
University of British Columbia, May 1-31, 2012
Leah Edelstein-Keshet

Wave-pinning

Essential features

The behaviour



“Wave-pinning”

What is the underlying mechanism?



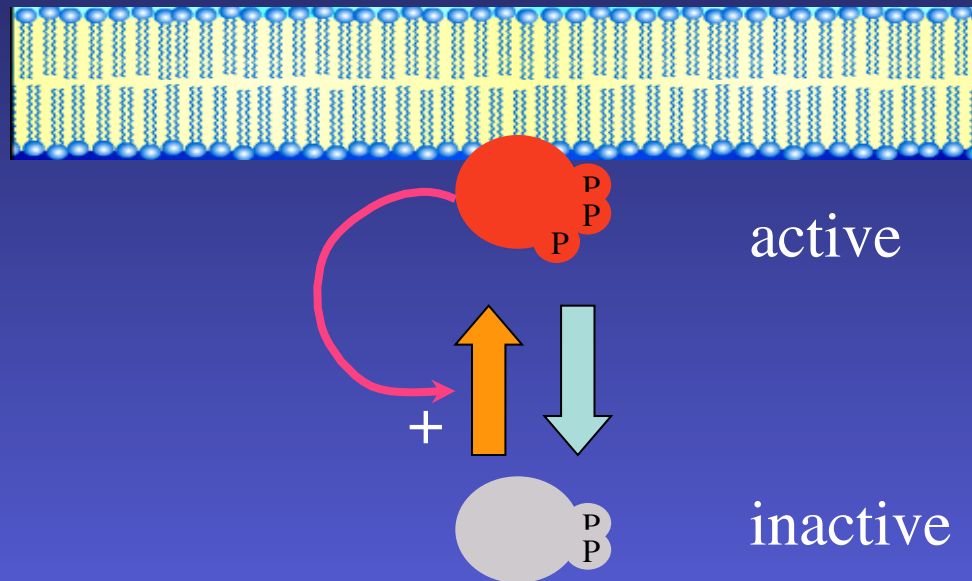
A Jilkiné



Y Mori

Mori Y, Jilkiné A, E-K L (2008) Biophysical Journal, 94: 3684-3697.

To investigate this, we study
simple system:



$$D_{\text{active}} \ll D_{\text{inactive}}$$

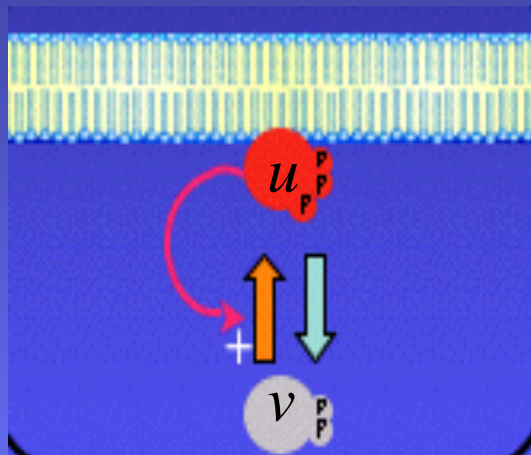
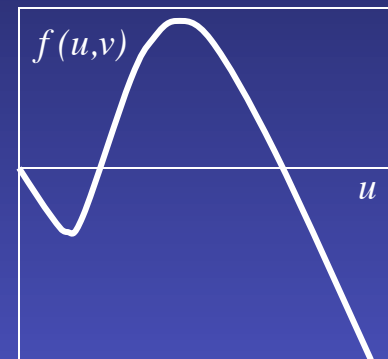
Mathematically:

Reaction-diffusion eqns with positive feedback terms

$$\begin{aligned}\frac{\partial u}{\partial t} &= D_u \frac{\partial^2 u}{\partial x^2} + f(u, v), \\ \frac{\partial v}{\partial t} &= D_v \frac{\partial^2 v}{\partial x^2} - f(u, v),\end{aligned}$$

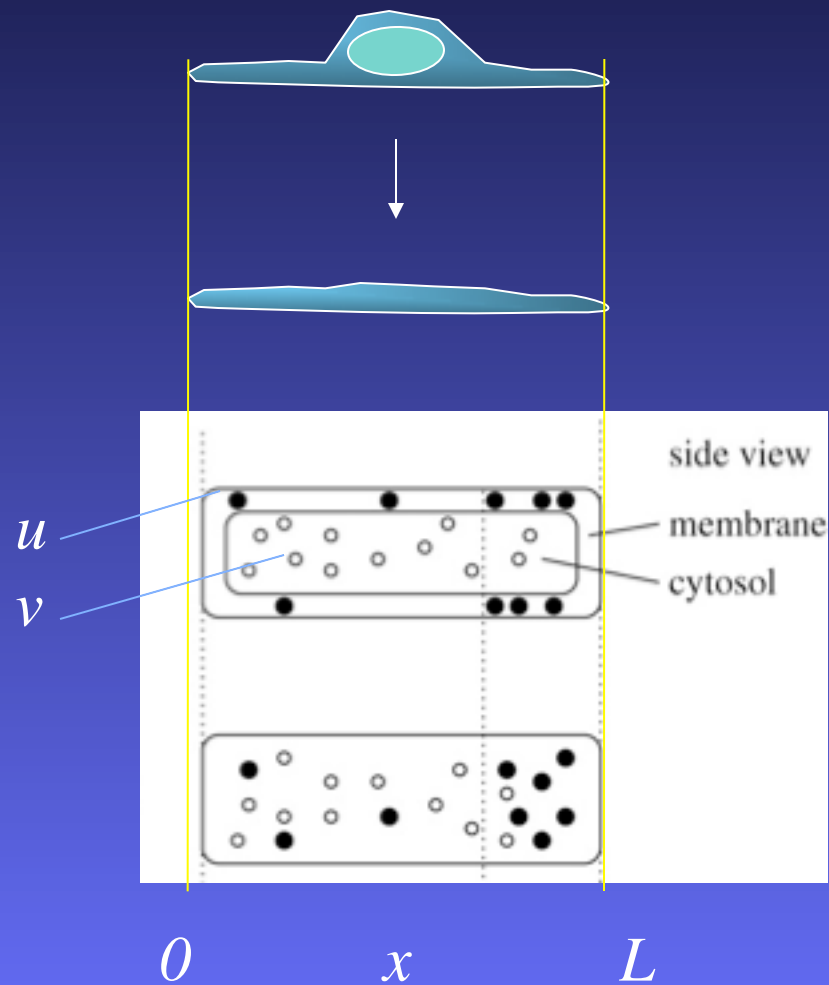
$$\int u + v = \text{Constant}$$

$$f(u, v) = \eta \left(\delta + \frac{\gamma u^2}{m^2 + u^2} \right) v - \eta u$$



$$D_u \ll D_v$$

Simplified 1D geometry:



$$D_u \ll D_v$$

Reduced model

Active form

Inactive form

$$\begin{aligned}\frac{\partial u}{\partial t} &= D_u \frac{\partial^2 u}{\partial x^2} + f(u, v), \\ \frac{\partial v}{\partial t} &= D_v \frac{\partial^2 v}{\partial x^2} - f(u, v),\end{aligned}$$

$$D_u \ll D_v$$

Typical kinetics

$$f(u, v) = \eta \left(\delta + \frac{\gamma u^2}{m^2 + u^2} \right) v - \eta u$$

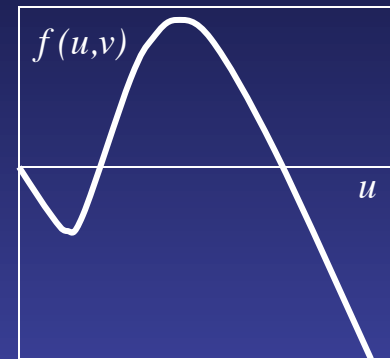
Neumann BCs

$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial x} = 0, \quad x = 0, L.$$



Moving fronts

$$\frac{\partial u}{\partial t} = D_u \frac{\partial^2 u}{\partial x^2} + f(u, v)$$

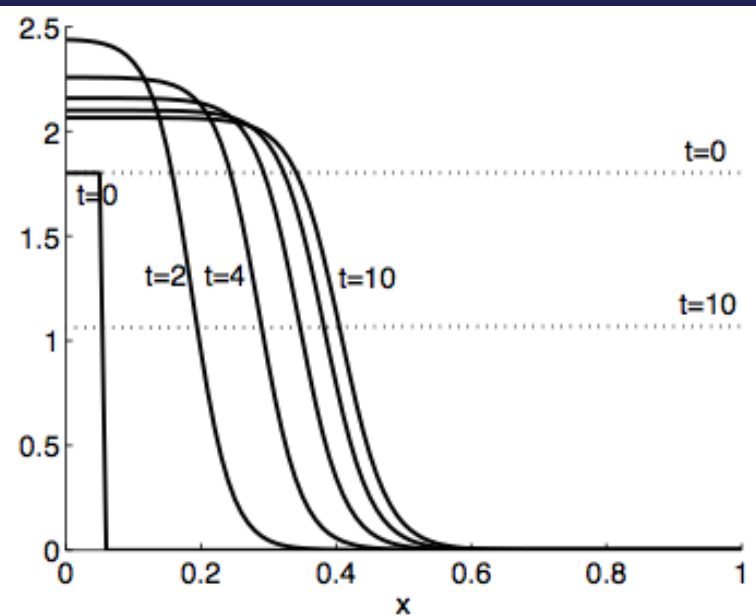
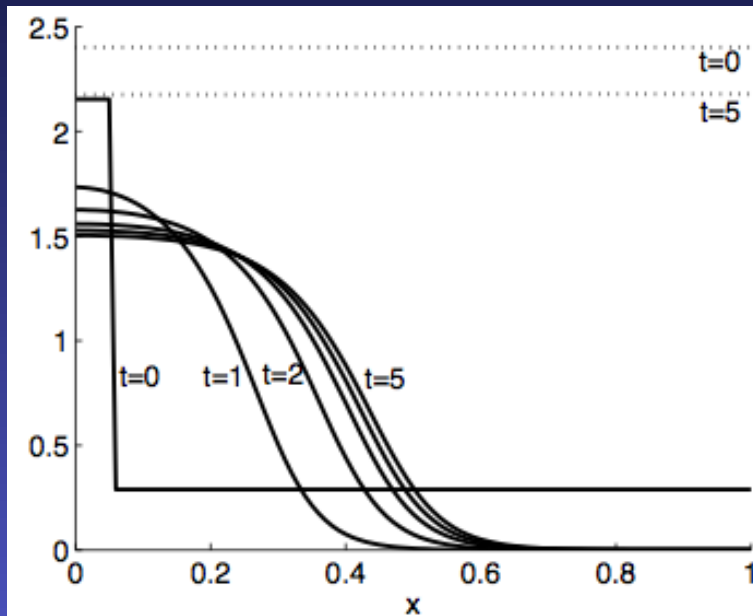


If v is fixed, and u eqn is on its own:

RD eqn with bistable kinetics.

This is known to have traveling wave solns
(moving fronts).

When coupled with the eqn for v ,
get a wave pinning phenomenon:



$$f(u, v) = \left(\delta + \frac{\gamma u^2}{1 + u^2} \right) v - u.$$

$$f(u, v) = u(1 - u)(u - 1 - v)$$

Rescaled (there is a small parameter)

$$\epsilon \frac{\partial u}{\partial t} = \epsilon^2 \frac{\partial^2 u}{\partial x^2} + f(u, v),$$

$$\epsilon \frac{\partial v}{\partial t} = D \frac{\partial^2 v}{\partial x^2} - f(u, v),$$

$$f(u, v) = \left(\delta + \frac{\gamma u^2}{1 + u^2} \right) v - u.$$

$$D_u \ll D_v$$

$$\epsilon^2 = \frac{D_u}{\eta L^2}, \quad D = \frac{D_v}{\eta L^2}$$

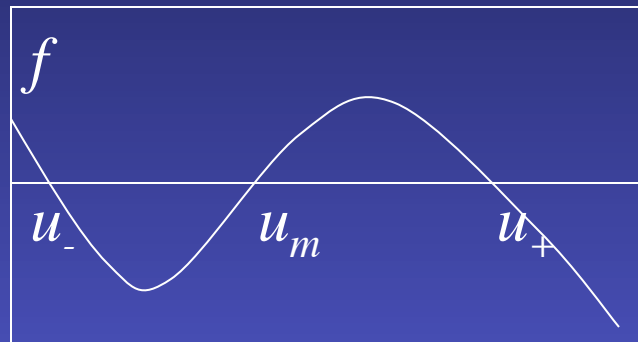
$$\int_0^1 (u + v) dx = K.$$

Conditions for Wave-pinning:

1. For v fixed in some range, $v_{min} < v < v_{max}$,

$f(u,v)=0$ has 3 roots

Shape of f :



1. There is a v_c in the above range such that
2. Conservation of $u+v$

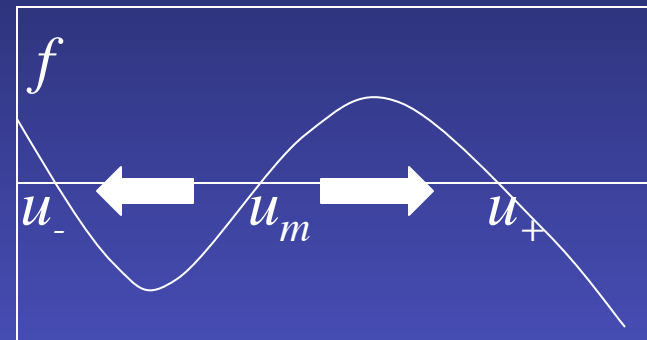
$$\int_{u_-}^{u_+} f(u, v_c) du = 0$$

Short time scale:

$$t_s = t/\epsilon$$

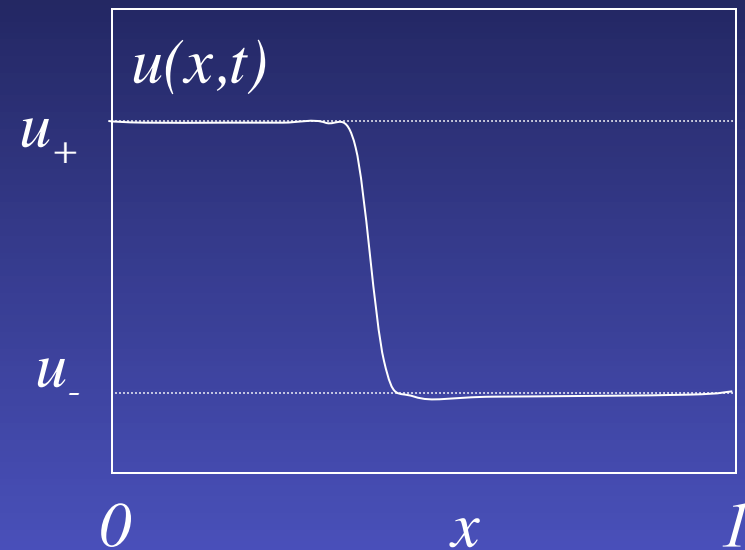
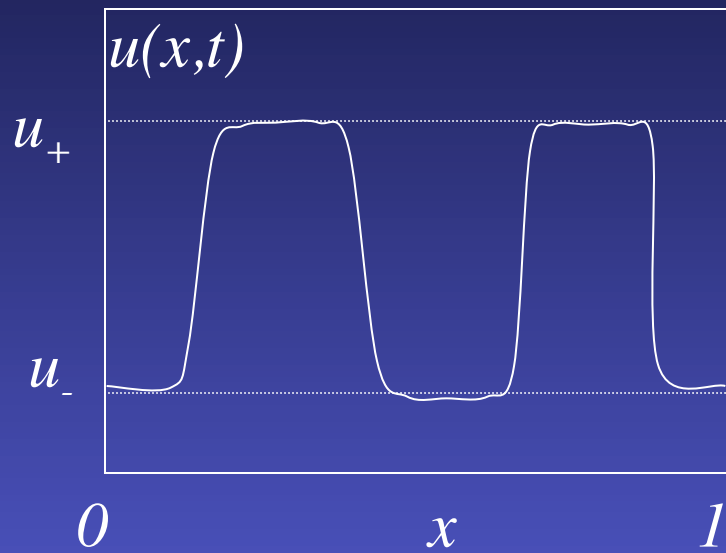
$$u = u_0 + \epsilon u_1, v = v_0 + \epsilon v_1$$

$$\begin{aligned}\frac{\partial u_0}{\partial t_s} &= f(u_0, v_0), \\ \frac{\partial v_0}{\partial t_s} &= D \frac{\partial^2 v_0}{\partial x^2} - f(u_0, v_0).\end{aligned}$$



To leading order, u_0 evolves to u_- or u_+
And v diffuses fast and attains uniform profile

Transition layers form on short timescale



Focus on single front solution (polarized domain)

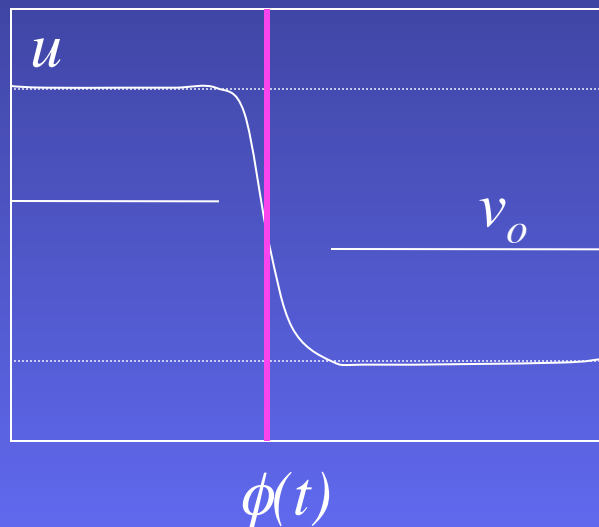
Intermediate time

Outer solution

$$0 = f(u_0, v_0),$$
$$0 = D \frac{\partial^2 v_0}{\partial x^2} - f(u_0, v_0).$$

Away from front $u = u_-$ or u_+

$$0 = D \frac{\partial^2 v_0}{\partial x^2}$$



With BCs on $[0,1]$, find that v_0 is constant left and at right of front

Intermediate time scale

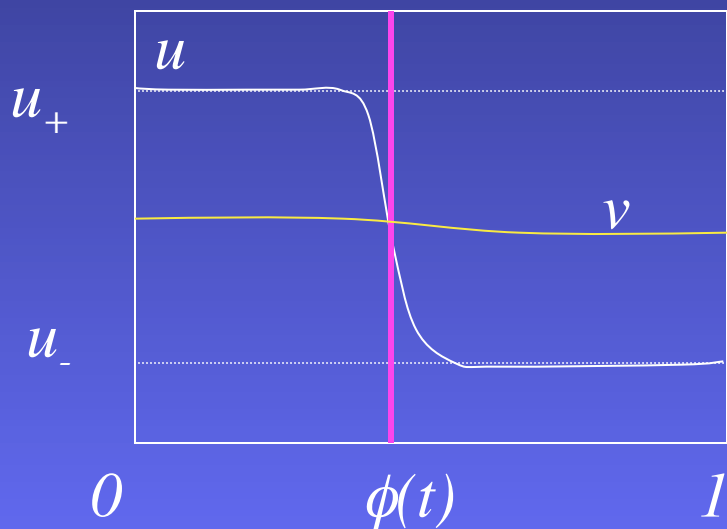
Stretched coordinate at the transition layer:

$$\xi = \frac{w}{\epsilon} = \frac{x - \phi(t)}{\epsilon}$$

Inner solution
to leading order:

$$\frac{\partial^2 U_0}{\partial \xi^2} - \frac{d\phi_0}{dt} \frac{\partial U_0}{\partial \xi} + f(U_0, V_0) = 0,$$

$$\frac{\partial^2 V_0}{\partial \xi^2} = 0.$$



V_0 becomes spatially uniform.

Matching inner and outer sols:

$$V_0 = v_0$$

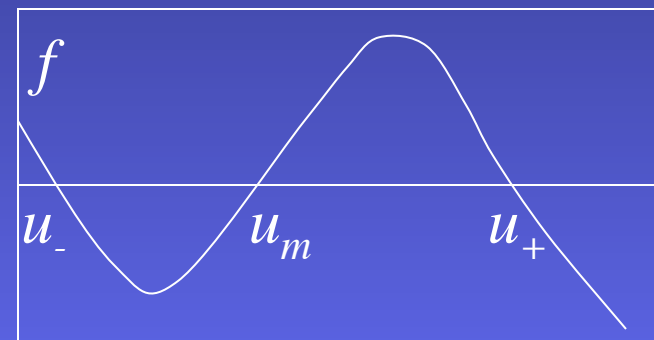
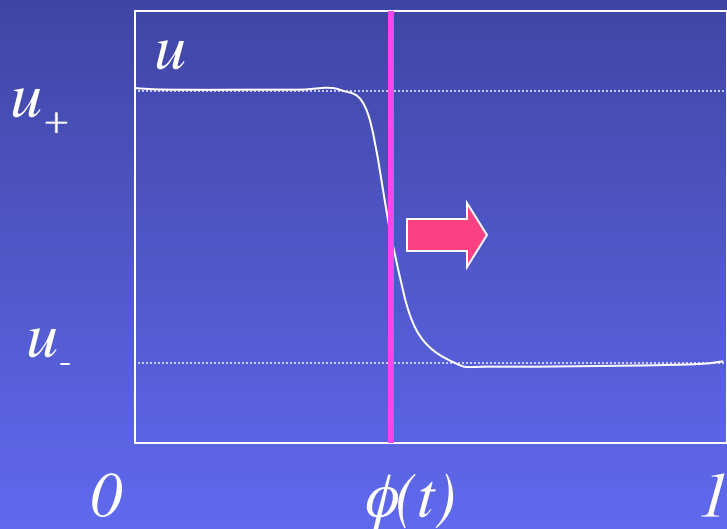
Front speed

In inner layer, V_0 is constant in space, so eqn for U_0 depends only on U_0 with V_0 a time-varying parameter.

We are back to bistable scalar RD eqn.

Speed of the wave

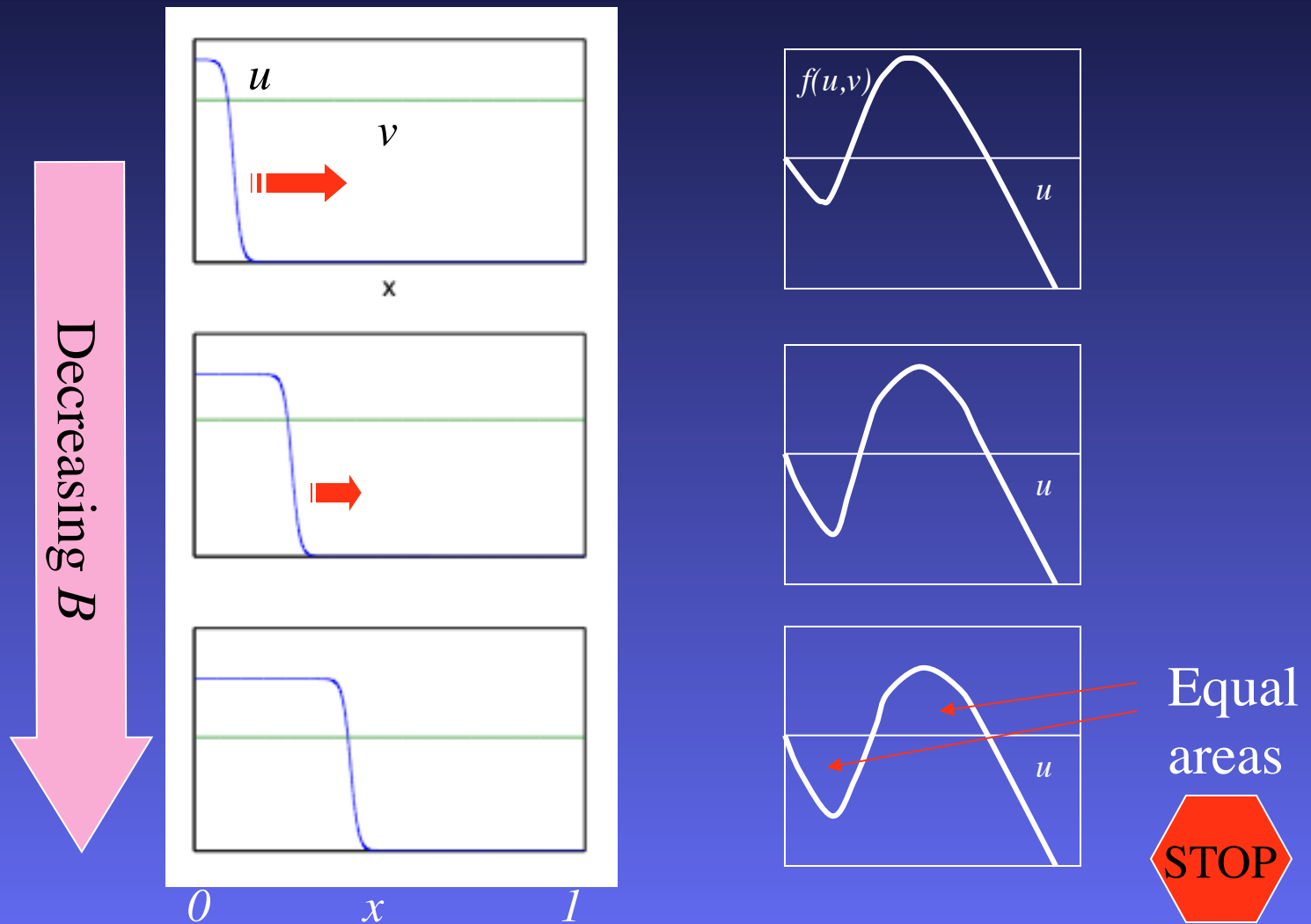
$$\frac{d\phi_0}{dt} \equiv c(V_0) = \frac{\int_{u_-(V_0)}^{u_+(V_0)} f(s, V_0) ds}{\int_{-\infty}^{\infty} \left(\partial U_0^\phi(\xi, V_0) / \partial \xi \right)^2 d\xi}.$$



Front motion depletes v

Front speed $d\phi_0/dt$, and dv_o/dt have opposite signs

Stalling of the wave

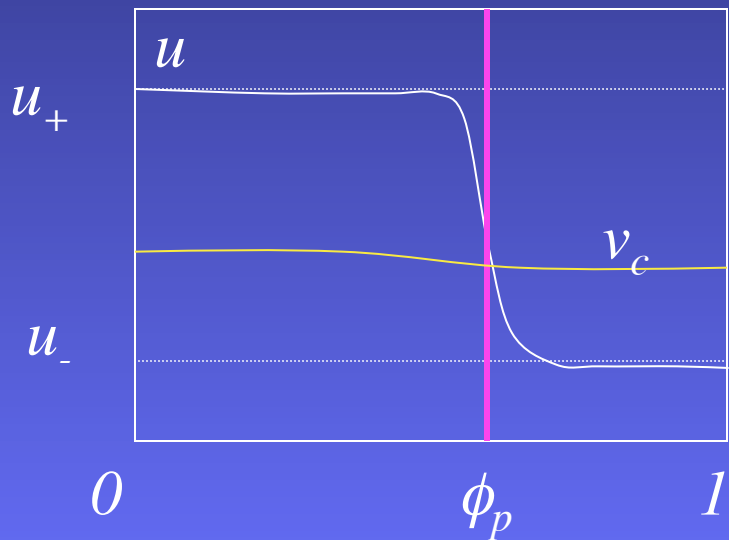
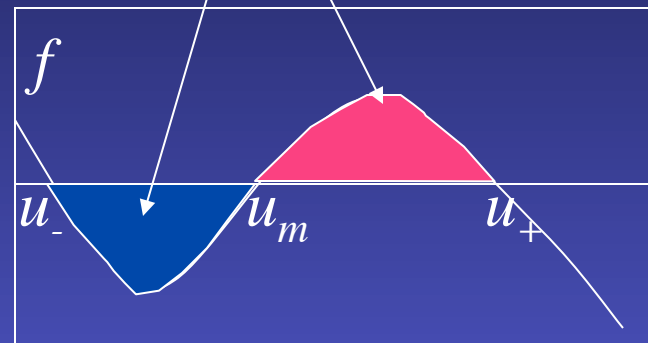


Wave stalls when

$$\int_{u_-}^{u_+} f(u, v) du = 0$$

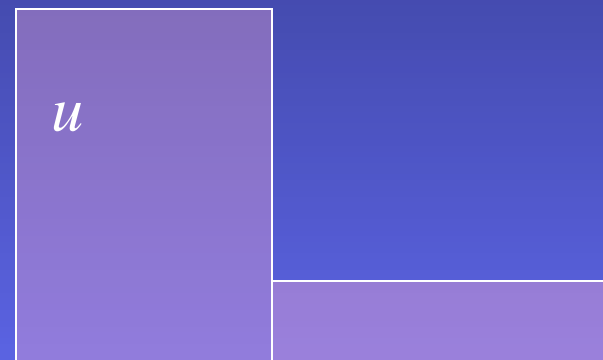
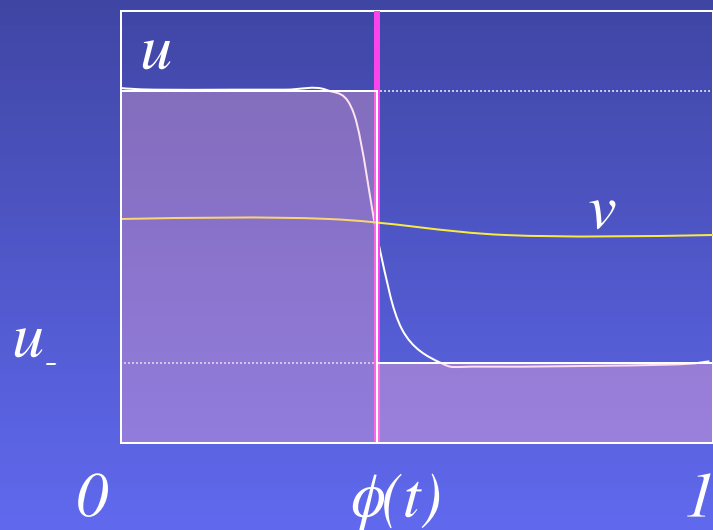
$$\frac{d\phi_0}{dt} \equiv c(V_0) = \frac{\int_{u_-(V_0)}^{u_+(V_0)} f(s, V_0) ds}{\int_{-\infty}^{\infty} (\partial U_0^\phi(\xi, V_0) / \partial \xi)^2 d\xi} = 0$$

(equal areas)



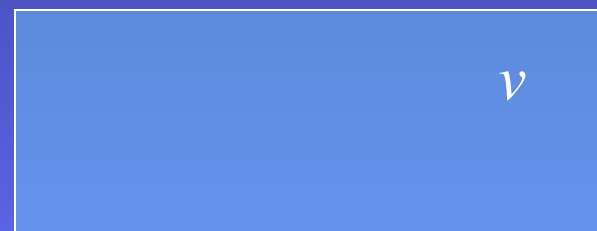
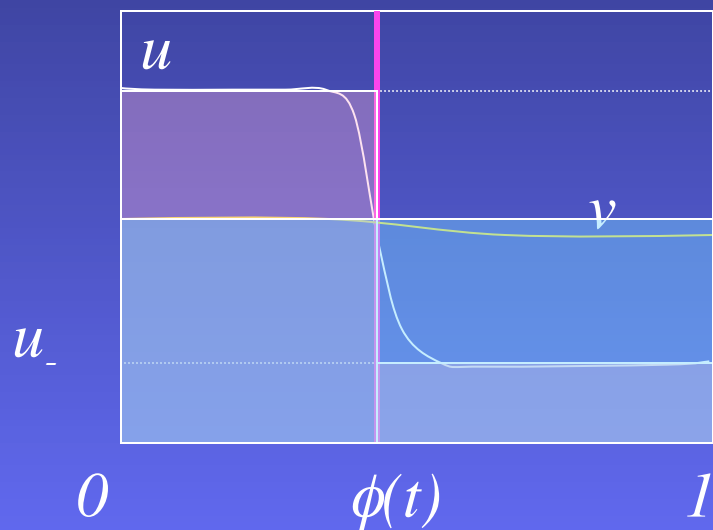
Where does the wave stop?

$$\int_0^1 (u + v) dx = K$$



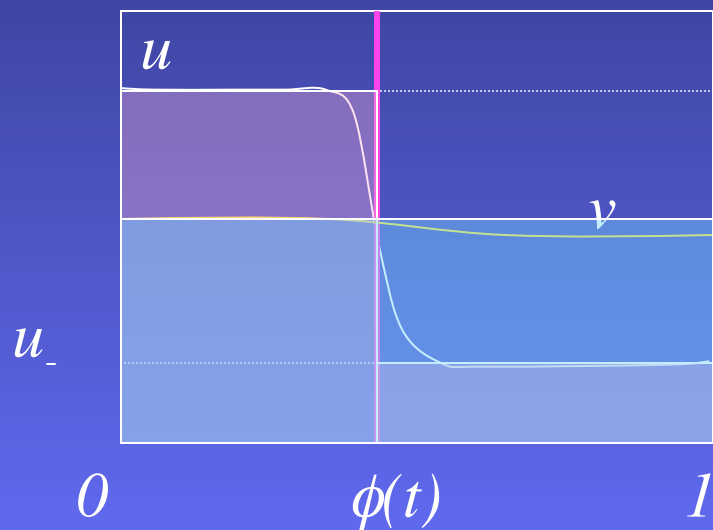
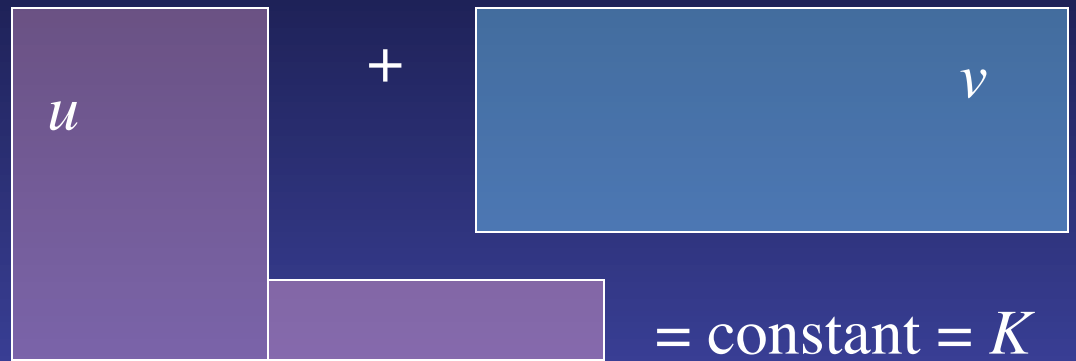
Where does the wave stop?

$$\int_0^1 (u + v) dx = K$$



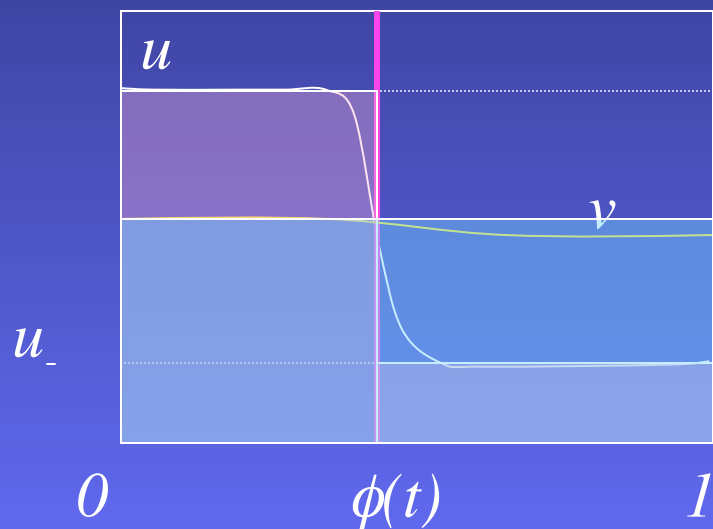
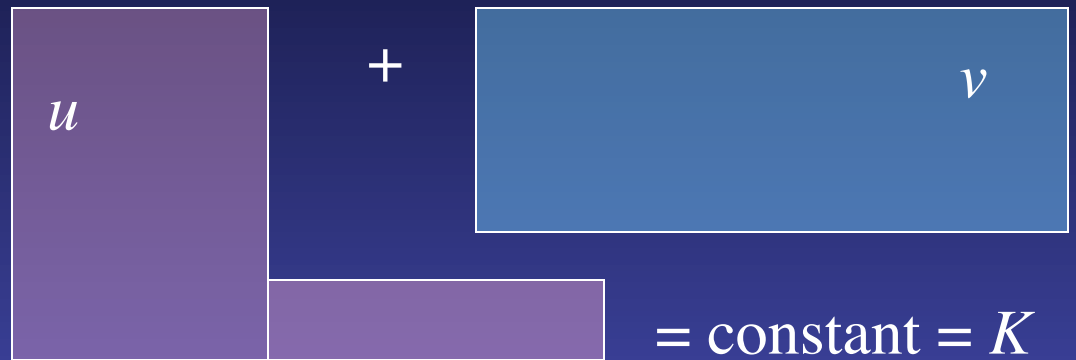
Conservation

$$\int_0^1 (u + v) dx = K$$



Conservation results in wave-pinning

$$\int_0^1 (u + v) dx = K$$



Wave stops when

$$v_c = K - u_+(v_c)\phi_p - u_-(v_c)(1 - \phi_p).$$

The mechanism:

Wave pinning results from the depletion of the inactive form and the fact that conservation of total amount of material is enforced.

Example: Cubic kinetics

$$f(u, v) = -(u - u_+(v))(u - u_m(v))(u - u_-(v))$$

Speed can be explicitly computed:

$$\frac{d\phi_0}{dt} \equiv c(V_0) = \frac{\int_{u_-(V_0)}^{u_+(V_0)} f(s, V_0) ds}{\int_{-\infty}^{\infty} (\partial U_0^\phi(\xi, V_0) / \partial \xi)^2 d\xi}.$$

Result:

$$c(v) = \frac{1}{\sqrt{2}} (u_+(v) - 2u_m(v) + u_-(v)).$$

Dynamics of front and stall position:

$$f(u, v) = u(1 - u)(u - 1 - v)$$

Speed of the wave :

$$\frac{d\phi_0}{dt} = \frac{1}{\sqrt{2}} \left(\frac{K - \phi}{1 + \phi} - 1 \right), \quad v_0 = \frac{K - \phi_0}{1 + \phi_0}$$

Stall position:

$$\phi_p = \frac{K - 1}{2}$$

Restriction for WP to occur:

$$\phi_p = \frac{K-1}{2}.$$

For wave to stall inside the domain: $0 < \phi_p < 1$.

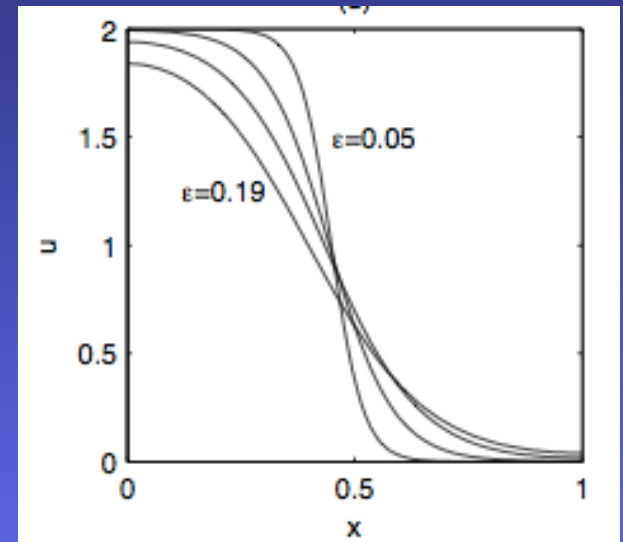
This can only happen if: $1 < K < 3$.

How does wave-pinning depend on parameters ?

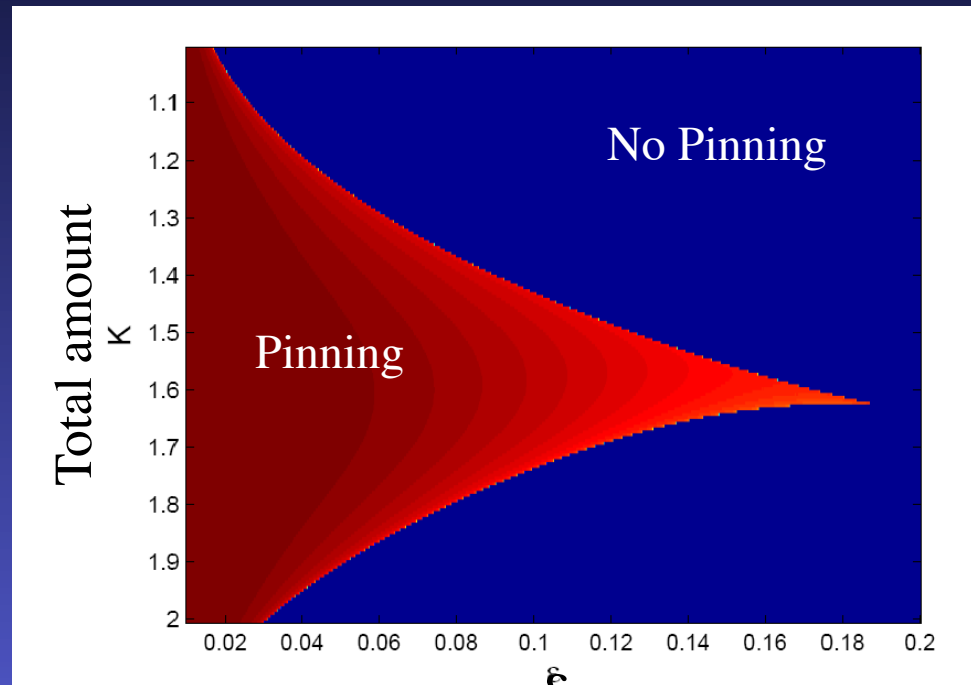
$$\epsilon^2 = \frac{D_u}{\eta L^2}, \quad D = \frac{D_v}{\eta L^2}$$

$$\int_0^1 (u + v) dx = K$$

- For $D > 0$, $1 < K < 3$, and ϵ sufficiently small, there is a stable front solution. As ϵ increases, shape and stall position change.
- For $\epsilon > \epsilon_c(D, K)$ this solution no longer exists.



wave-pinning loss

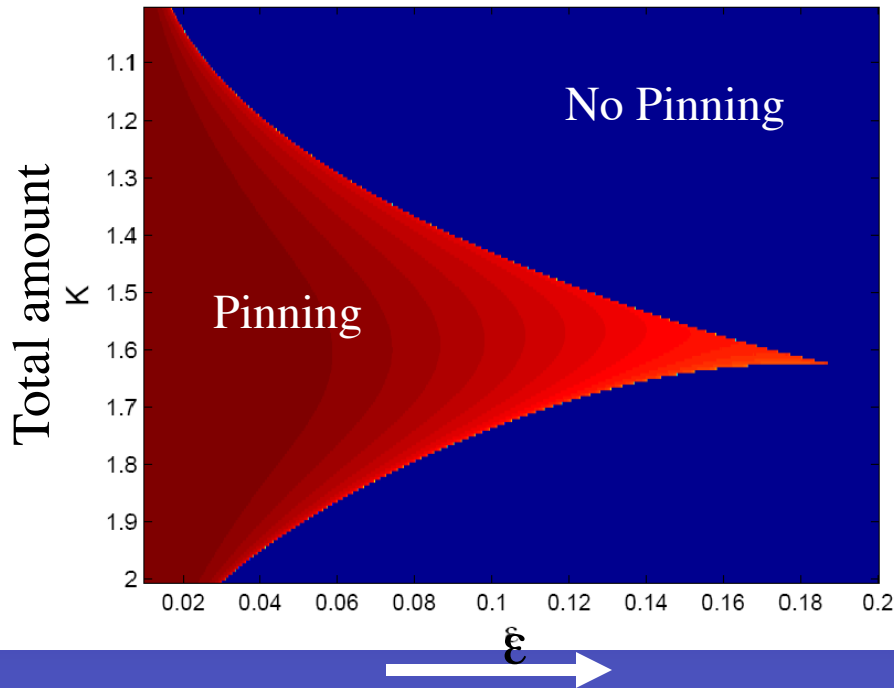


Diffusion D_u , \uparrow

Domain size L \downarrow

reaction rate \downarrow

Implications



If membrane diffusion too fast, or domain too small or reaction rate too slow, no wave-pinning so cell can't polarize.

Diffusion D_u , \uparrow
Domain size L , \downarrow
reaction rate \downarrow