

Travelling waves and Wave-Pinning in a set of PDEs.

Goals: - To explore travelling wave speed in a single Reaction-Diffusion (RD) equation

- To examine a pair of RD eqns with mass conservation and "cubic" type kinetics
- Based on vastly different rates of diffusion, to arrive at a scaled ^{dimensionless} model with a small parameter (ϵ)
- To exploit the small parameter in asymptotic expansions both near an interface of the wave ("inner soln") and far away ("outer soln")
- To arrive at an ODE that characterizes the motion of that interface, and to show that it can stall (\equiv wavepinning) under the appropriate conditions.

Preliminaries : Speed of travelling wave in a single RD eqn.

Let $u(x,t)$ be variable of interest on $-\infty \leq x \leq \infty$

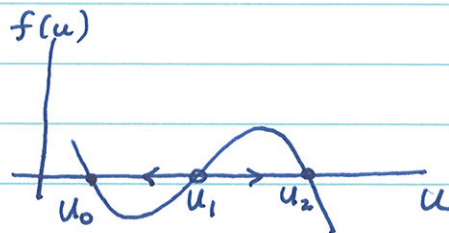
Suppose u satisfies

$$u_t = f(u) + D u_{xx}$$

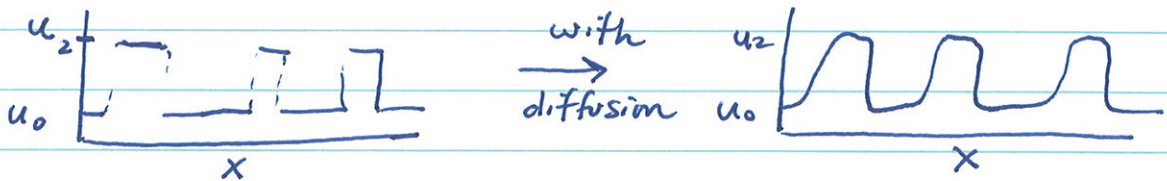
(reaction) (diffusion)

with no-flux boundary conditions (BC's) : $u_x = 0$
at $x = \pm \infty$

For example :
"cubic kinetics"

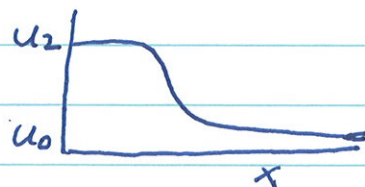


In the well-mixed system, such kinetics lead to one of two stable steady states, u_0 or u_2 (depending on initial conditions.) We would expect that spatially, we could obtain a solution to the pde that "mixes" these steady states:



Here, we have shown that diffusion smooths out the sharp transitions.

We will be interested here in single interfaces:

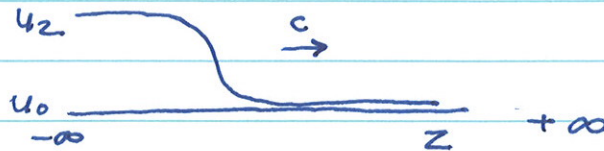


we ask whether such solns exist, and, in particular whether they are stationary or travelling waves. We are, moreover, interested in speed and direction of motion.

Travelling waves: To look for such soln's, we transform to a moving coordinate system (moving at speed c)

$$\text{Let } z = x - ct \quad U(z(x,t)) = u(x,t)$$

In this coordinate system, the wave will look like:



$$\text{Use chain rule to write } u_x = \frac{\partial u}{\partial x} = \frac{dU}{dz} \frac{\partial z}{\partial x} = \frac{dU}{dz} \cdot 1$$

$$u_t = \frac{\partial u}{\partial t} = \frac{dU}{dz} \frac{\partial z}{\partial t} = \frac{dU}{dz} (-c)$$

Transforming to the new variables leads to

$$-c \frac{dU}{dz} = f(U) + D \frac{d^2 U}{dz^2} \quad \text{and } \frac{dU}{dz} = 0 \text{ for } z = \pm\infty$$

Wave speed: First note that the following small calculation

is useful:
$$\frac{d}{dz} \left(\left(\frac{du}{dz} \right)^2 \right) = 2 \frac{du}{dz} \cdot \frac{d}{dz} \left(\frac{du}{dz} \right) = 2 \frac{du}{dz} \frac{d^2u}{dz^2}$$

Our eqn, in the moving coordinate frame is

$$-c \frac{du}{dz} = f(u) + D \frac{d^2u}{dz^2}$$

Multiply by $\frac{du}{dz}$:

$$\begin{aligned} -c \left(\frac{du}{dz} \right)^2 &= f(u) \frac{du}{dz} + D \frac{d^2u}{dz^2} \frac{du}{dz} \\ &= f(u) \frac{du}{dz} + \frac{D}{2} \frac{d}{dz} \left(\left(\frac{du}{dz} \right)^2 \right) \quad (\text{see above}) \end{aligned}$$

Integrate both sides:

$$\begin{aligned} -c \int_{-\infty}^{\infty} \left(\frac{du}{dz} \right)^2 dz &= \int_{-\infty}^{\infty} f(u) \frac{du}{dz} dz + \frac{D}{2} \int_{-\infty}^{\infty} \frac{d}{dz} \left(\left(\frac{du}{dz} \right)^2 \right) dz \\ &= \int_{u(-\infty)}^{u(\infty)} f(u) du + \frac{D}{2} \left(\frac{du}{dz} \right)^2 \Big|_{-\infty}^{\infty} \end{aligned}$$

BCs $\Rightarrow \frac{du}{dz} = 0$ at $z = \pm\infty$
so this term vanishes.

Since we look for a

moving front as shown on previous page, $u(-\infty) = u_2$
 $u(+\infty) = u_0$

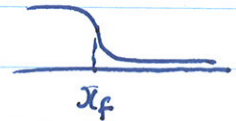
so

$$c = - \int_{u_2}^{u_0} f(u) du / \int_{-\infty}^{\infty} \left(\frac{du}{dz} \right)^2 dz$$

Remarks:

- Same calculation works equally well for non-constant front speed, e.g. set

$$z = x - x_f(t)$$



where $x_f(t)$ is the "position" of the moving front (e.g. some coordinate on that interface, such as inflection pt. or point corresponding to arbitrary u value).

$$\text{Then } z = z(x, t) \quad \frac{\partial z}{\partial x} = 1 \quad \frac{\partial z}{\partial t} = -\frac{dx_f}{dt} = -v(t)$$

where $v(t)$ is the time-dependent wave speed.

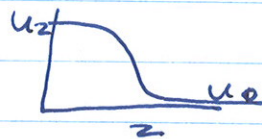
- The above calculation, ^{on previous page} may not provide an explicit value of the speed c , but we can extract from it the direction of motion ($c > 0$ or $c < 0$) and whether the wave can ever be stationary ($c = 0$). Note that the denominator is always > 0 .

Also: $c = 0 \Leftrightarrow$ stationary wave \Leftrightarrow

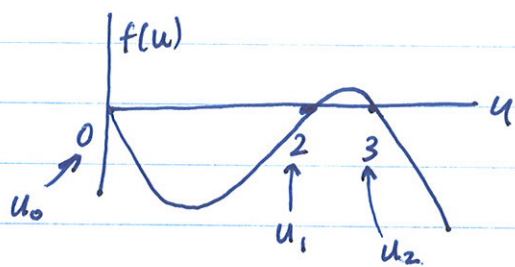
$$\int_{u_2}^{u_0} f(u) du = 0$$

Exercise: (1) Suppose $f(u) = (u-u_0)(u_1-u)(u-u_2)$

If $u_0 = 0$, $u_1 = 2$ and $u_2 = 3$, which way would a front solution of the form $u(z)$ move?



Soln: sketch $f(u)$; a cubic. Even a rough sketch will



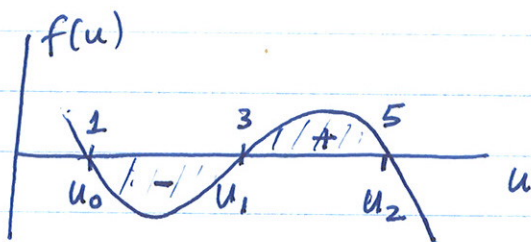
convince us that

$$c = - \int_{u_0}^{u_2} f(u) du = - \left[\text{positive quantity} \right]$$

$c < 0$ so move to left.

(2) For the same form of f , given that $u_0 = 1$ and $u_1 = 3$, for what value of u_2 would we expect a stationary wave?

Soln: To answer this, we look for u_2 such that the plot of $f(u)$ has 0 integral over the range (u_0, u_2)



By symmetry of the cubic, picking $u_2 = 5$ will do it.

since have "equal and opposite" "areas" (shown shaded)

thus

$$\int_{u_2}^{u_0} f(u) du = \int_{u_2}^{u_1} + \int_{u_1}^{u_0} = 0$$